

# Emergent magnetic order from weak-crystal field, Octupolar quantum spin ice, Hole doped $\text{Sr}_2\text{IrO}_4$

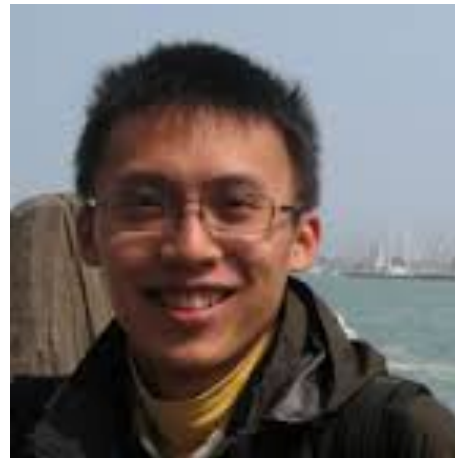
Gang Chen  
June 11, 2016

as usual psychology, one is more excited about the new work. today i will present some preliminary results of these new problems.



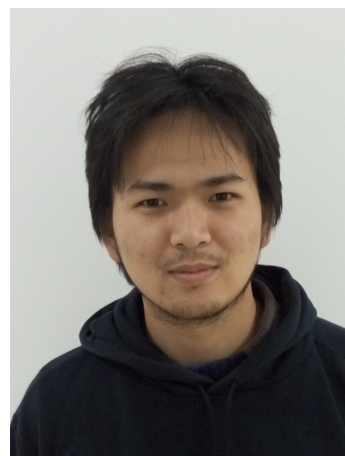
# Outline

- Emergent magnetic order due to “weak” crystal field gap.



Yi-Ping Huang  
(Univ of Colorado, Boulder)

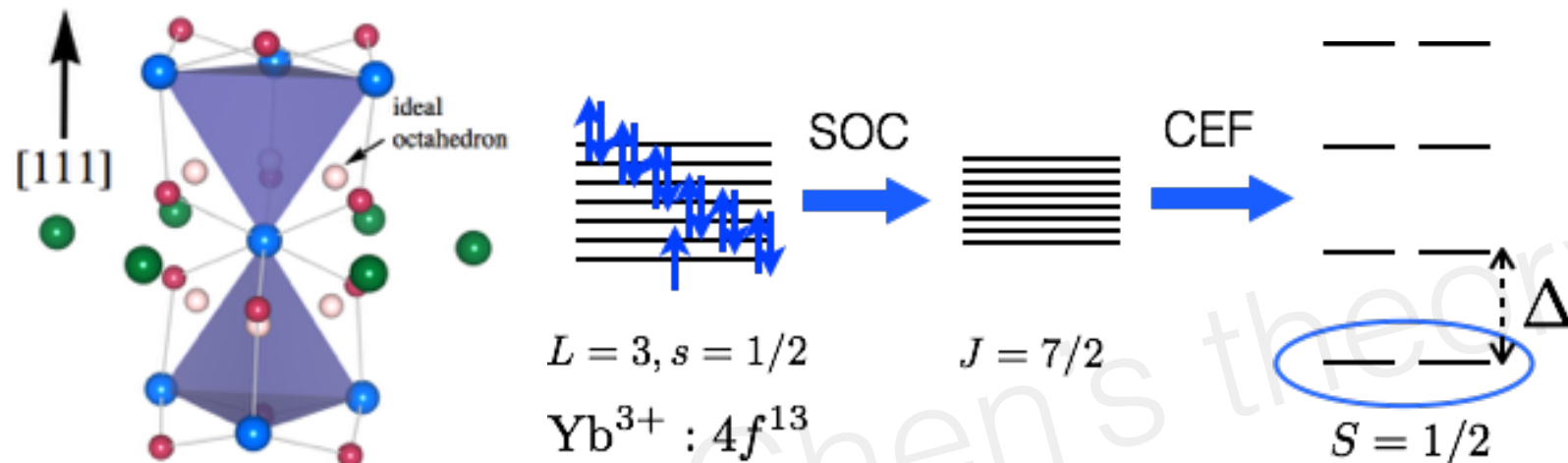
- Field driven Anderson-Higgs' transition in octupolar quantum spin ice.



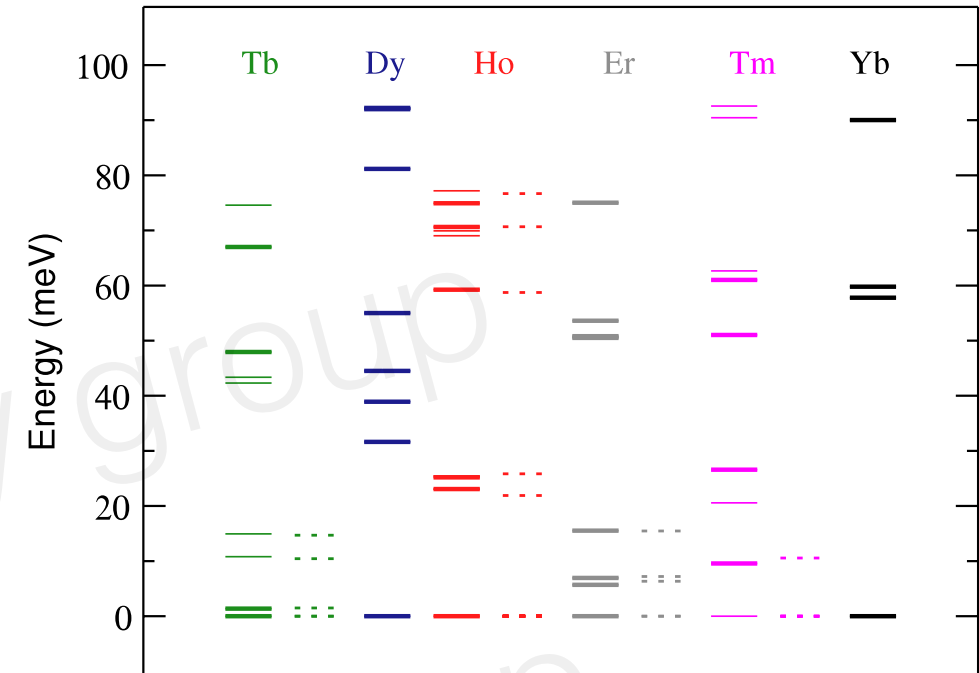
Yao-Dong Li  
(Dept of computer sciences, Fudan)

- Hole doped  $\text{Sr}_2\text{IrO}_4$ : the difference from cuprates ! (by myself)

# Crystal field energy levels



e.g.  $\text{YbMgGaO}_4$ , Yao-dong Li, XQ Wang, GC, 2016



**Figure 1.** The computed CEF energy scheme drawn for the R ions in the  $\text{R}_2\text{Ti}_2\text{O}_7$  pyrochlore series and comparison with experimental values when available, as extracted from inelastic neutron scattering measurements. Energy levels are given in units of millielectronvolts.

We often think that the ground state doublet controls the low temperature magnetic properties.

This happens when the crystal field gap is much larger than the temperature and exchange interaction scales.

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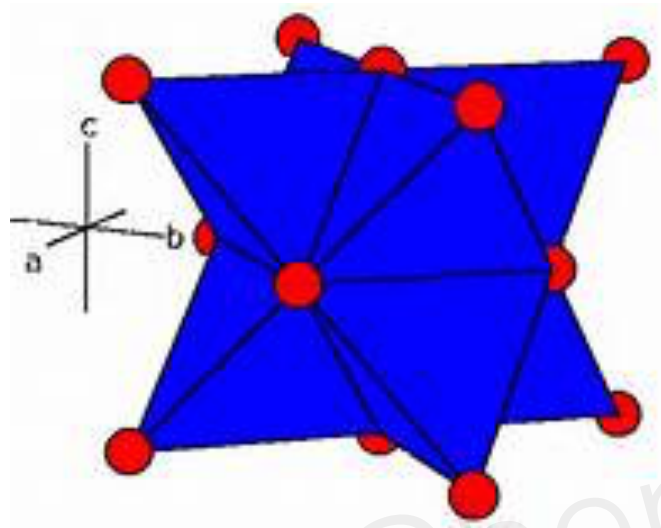
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Figure 1 consists of two panels. Panel (a) is a 2D color plot of phonon dispersion. The vertical axis is 'Energy transfer (meV)' ranging from 0 to 3. The horizontal axis is ' $(h,h,h)$  rlu' ranging from 0 to 2.5. A color bar at the top indicates intensity from 0 to 4. An inset shows a zoomed-in view of the CEF level region. Panel (b) is a plot of 'Phonon integrated intensity (a.u.)' vs. 'Energy transfer (meV)'. The vertical axis ranges from 0.0 to 1.0. The horizontal axis ranges from 0.0 to 3.0. Data points with error bars are shown, along with a solid line fit. A shaded gray region indicates the CEF level.

Question: how if the crystal field gap is “weak”?  
or equivalently, the exchange is strong.



# Model Calculation



FCC lattice

⋮ ⋮ ⋮

$$\begin{array}{cc} S^z = 1 & S^z = -1 \\ \hline & \\ S^z = 0 & \end{array}$$

$$\begin{aligned} \tau^z &= \frac{1}{2} s^z \\ \tau_i^\pm &= \frac{1}{2} (s_i^\pm)^2 \end{aligned}$$

$$H_{ex} = \frac{1}{2} \sum_{\langle ij \rangle} J_{zz} \tau_i^z \tau_j^z - J_{\pm} (\tau_i^+ \tau_j^- + h.c.)$$



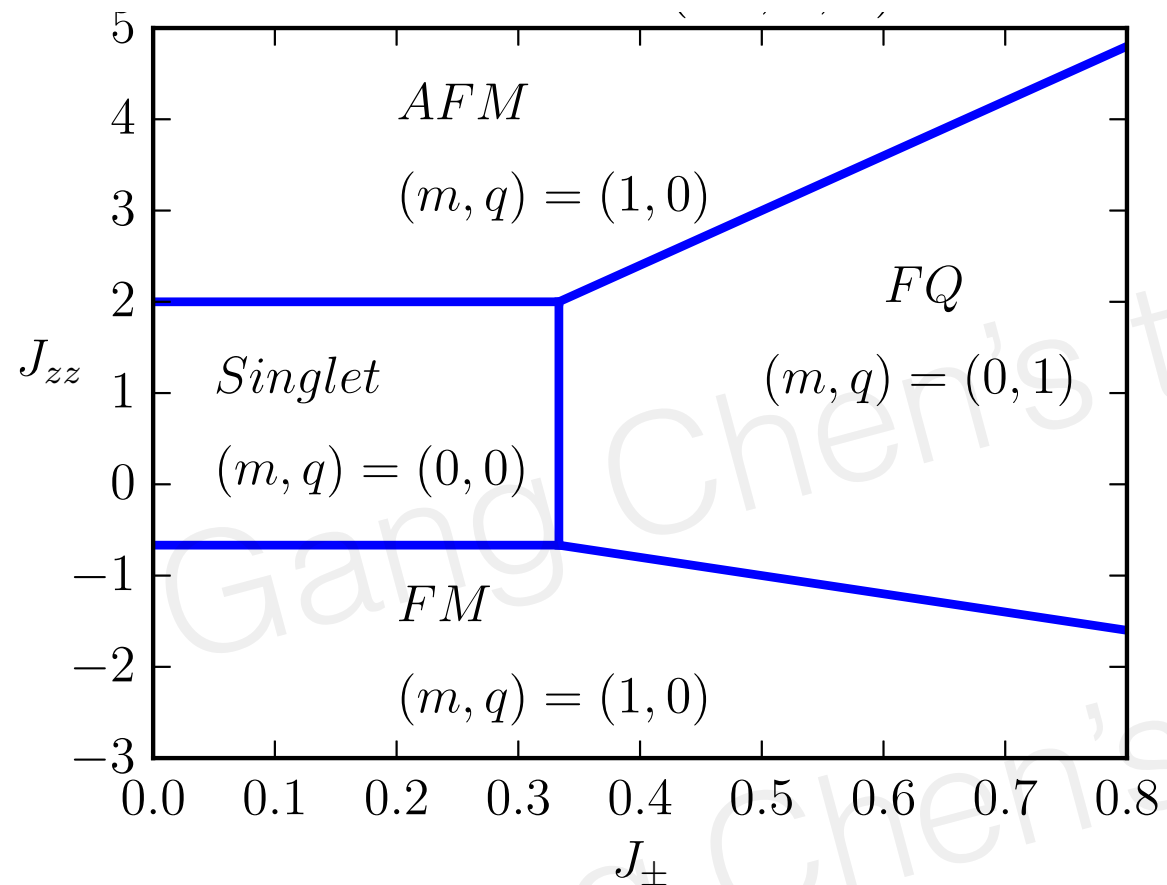
$$H =$$

Since this observation is general, we can actually demonstrate this physics can happen in any other lattice.

Our formalism here can be easily adapted to more general situations.

$$H_{CEF}$$

# Phase diagram



m is the magnetic dipolar order  
q is the magnetic quadrupolar order

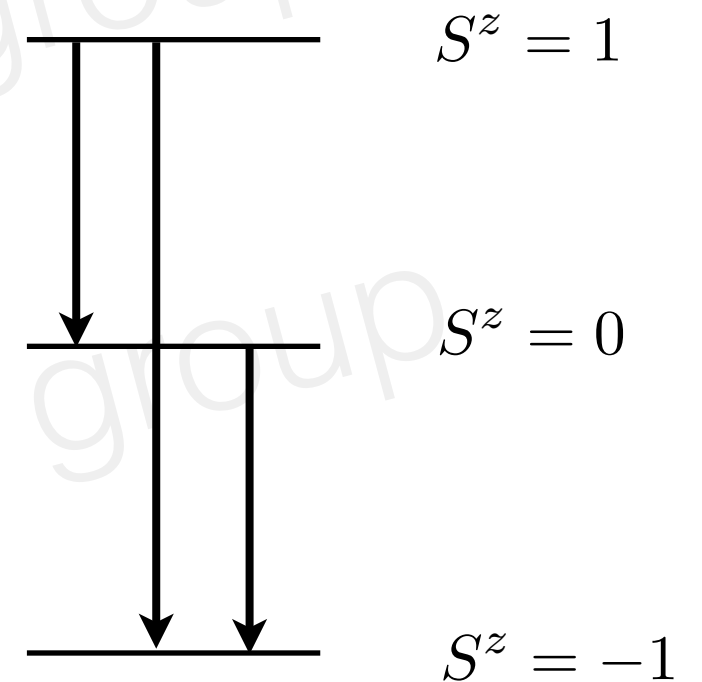
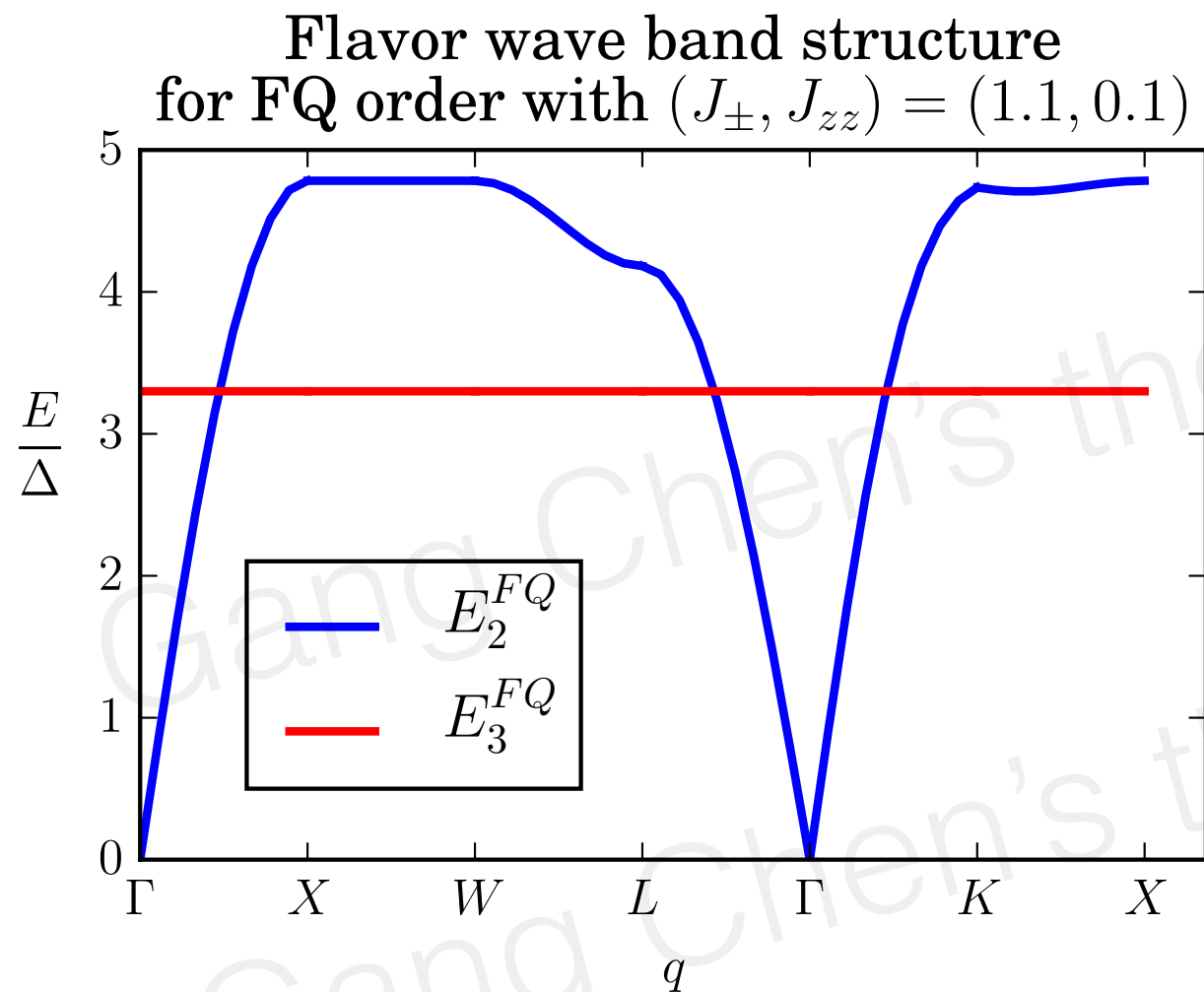
FQ = Ferroquadrupolar order  
or Ferro-Spin-Nematic order

$$\langle (s_i^+)^2 \rangle \neq 0$$

$\Delta$  is set as the energy unit

two spins are condensed,  
one spin is not condensed.

# Excitation spectrum

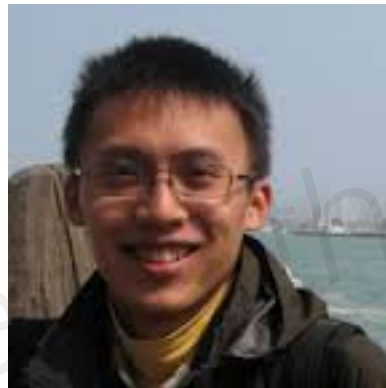


excitation of the ferroquadrupolar phase

$$- J_{\pm} [(s_i^+)^2 (s_j^-)^2 + h.c.]$$

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Yi-Ping Huang  
(CU-Boulder)

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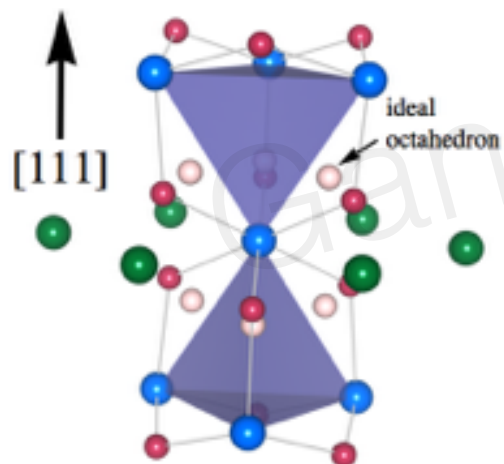
Yao-Dong Li  
(Fudan)

- Hole doped  $\text{Sr}_2\text{IrO}_4$ : the difference from cuprates !

# Dipole-octupole doublet

- Why is this Kramers doublet so special?

ONE-dimensional representations of the point group



$$R(2\pi/3)|J^z = \pm 3/2\rangle = -|J^z = \pm 3/2\rangle$$

$$R(2\pi/3) \equiv e^{-i\frac{2\pi}{3}J^z} = e^{-i\frac{2\pi}{3} \times (\pm\frac{3}{2})} = e^{\mp i\pi} = -1$$

$$|J^z = +3/2\rangle \xrightarrow{\text{time reversal}} |J^z = -3/2\rangle$$

why special,

they are one dimensional rep of po group.

in particular, if you look at the 3-fold rotation operation, under this rotation each state stay invariant, except a n sign. they do not transform into each other under the point group transformation.

simple algebra

this is very different from 2-dim irep where these two states of the doublet are mixed.

These two state are degenerate under time reversal, deg protected by time reversal.



# More generally, ...

- Also applies to 4f electron moments on pyrochlore

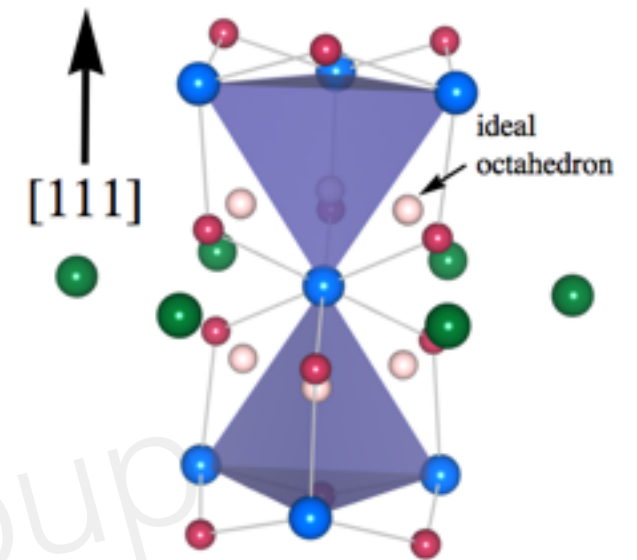
$$J = \frac{3}{2}, \frac{9}{2}, \frac{15}{2}, \dots$$

with the local crystal field Hamiltonian

$$H_{\text{cf}} = 3B_2^0(J^z)^2 + \dots \quad \text{if } B_2^0 < 0$$

e.g. local doublet wavefunction of  $\text{Dy}^{3+}$  ( $J = \frac{15}{2}$ ) in  $\text{Dy}_2\text{Ti}_2\text{O}_7$

$$|\phi_0^\pm\rangle = 0.981|\pm\frac{15}{2}\rangle \pm 0.190|\pm\frac{9}{2}\rangle - 0.022|\pm\frac{3}{2}\rangle \mp 0.037|\mp\frac{3}{2}\rangle + 0.005|\mp\frac{9}{2}\rangle \pm 0.001|\mp\frac{15}{2}\rangle$$

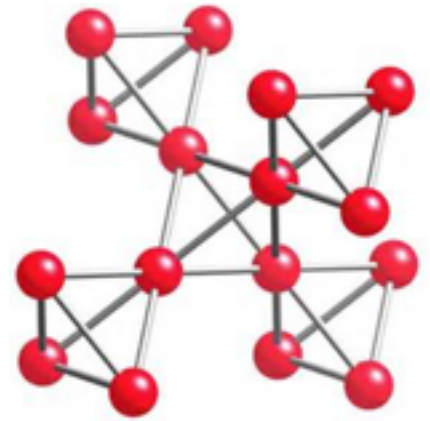


more generally, it applies to xxxx.  
if the local crystal field hamiltnian is  
easy axis like.

Emphasis: what matters is the wavefunction, not the spin value !

- May generally apply to any Kramers' doublets with  $J > 1/2$  !

e.g, Ce: **Ce<sub>2</sub>Sn<sub>2</sub>O<sub>7</sub>**



PRL **115**, 097202 (2015)

PHYSICAL REVIEW LETTERS

week ending  
28 AUGUST 2015

### Candidate Quantum Spin Liquid in the Ce<sup>3+</sup> Pyrochlore Stannate Ce<sub>2</sub>Sn<sub>2</sub>O<sub>7</sub>

Romain Sibille,<sup>1,\*</sup> Elsa Lhotel,<sup>2</sup> Vladimir Pomjakushin,<sup>3</sup> Chris Baines,<sup>4</sup> Tom Fennell,<sup>3,†</sup> and Michel Kenzelmann<sup>1</sup>

$4f^1$  ion in  $D_{3d}$  local symmetry to the susceptibility was realized between  $T = 1.8$  and 370 K, and the resulting calculation of the single ion magnetic moment is shown in Fig. 2(c). The wave functions of the ground state Kramers doublet correspond to a linear combination of  $m_J = \pm 3/2$  states. The fitted coefficients result in energy levels at  $50 \pm$

$$\text{Ce}^{3+} (4f^1, {}^2F)$$
$$J = \frac{5}{2}$$

the wavefunction should be a combination of states with odd integer values of  $J$ .

this really because of the fact that they are one-dimensional representations, even in the YbTiO<sub>2</sub>, such as the excited doublets.

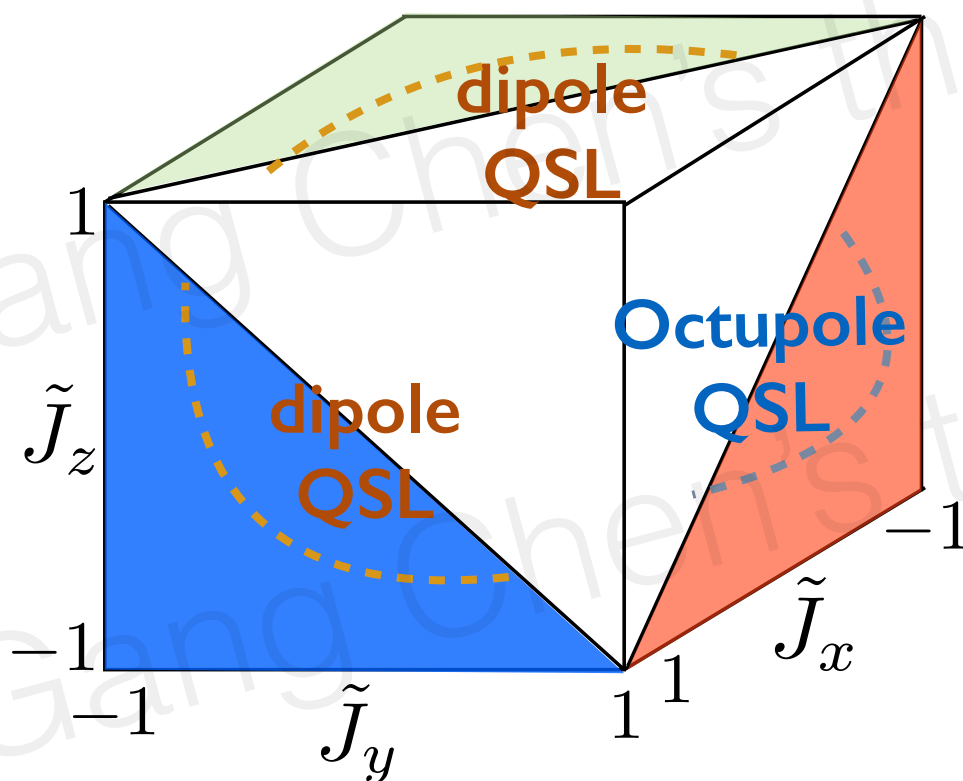
experimentalists of this paper do not know what they really talk about.

But this sentence means a lot to us. It means that gs doublet is a DO doublet, and the model is described by XX

XYZ model is the generic model that describes the interaction between DO doublets.

$$H_{\text{XYZ}} = \sum_{\langle ij \rangle} \mathcal{J}_x \tau_i^x \tau_j^x + \mathcal{J}_y \tau_i^y \tau_j^y + \mathcal{J}_z \tau_i^z \tau_j^z$$

unlike XXZ model, XYZ model is richer



3D phase diagram

Each component (not just  $S_z$ ) can be emergent electric field, depending on the parameters !

Study phase on a cube:  $-1 \leq \tilde{J}_{x,y,z} \leq 1$ .

# Emergent Quantum Electrodynamics of U(1) QSL

as quantum spin ice is  
there is no long range  
a new phase of matter  
in the Landau's paradigm

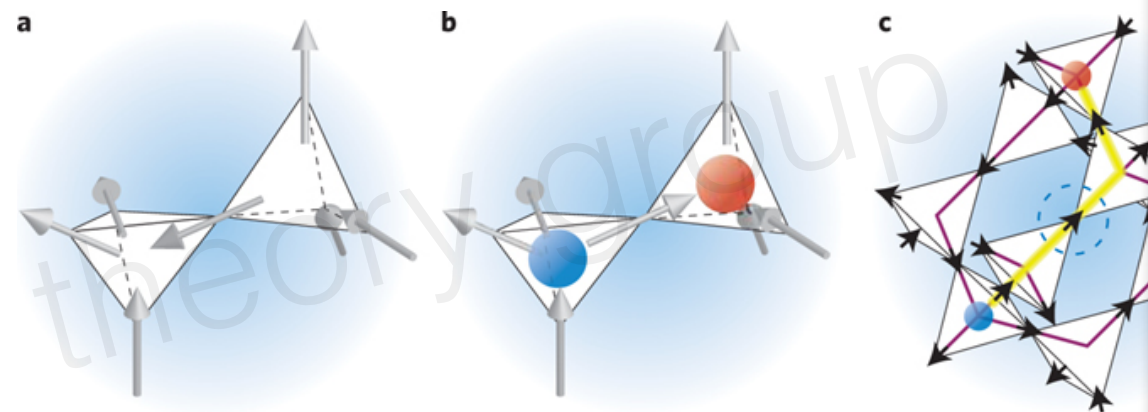
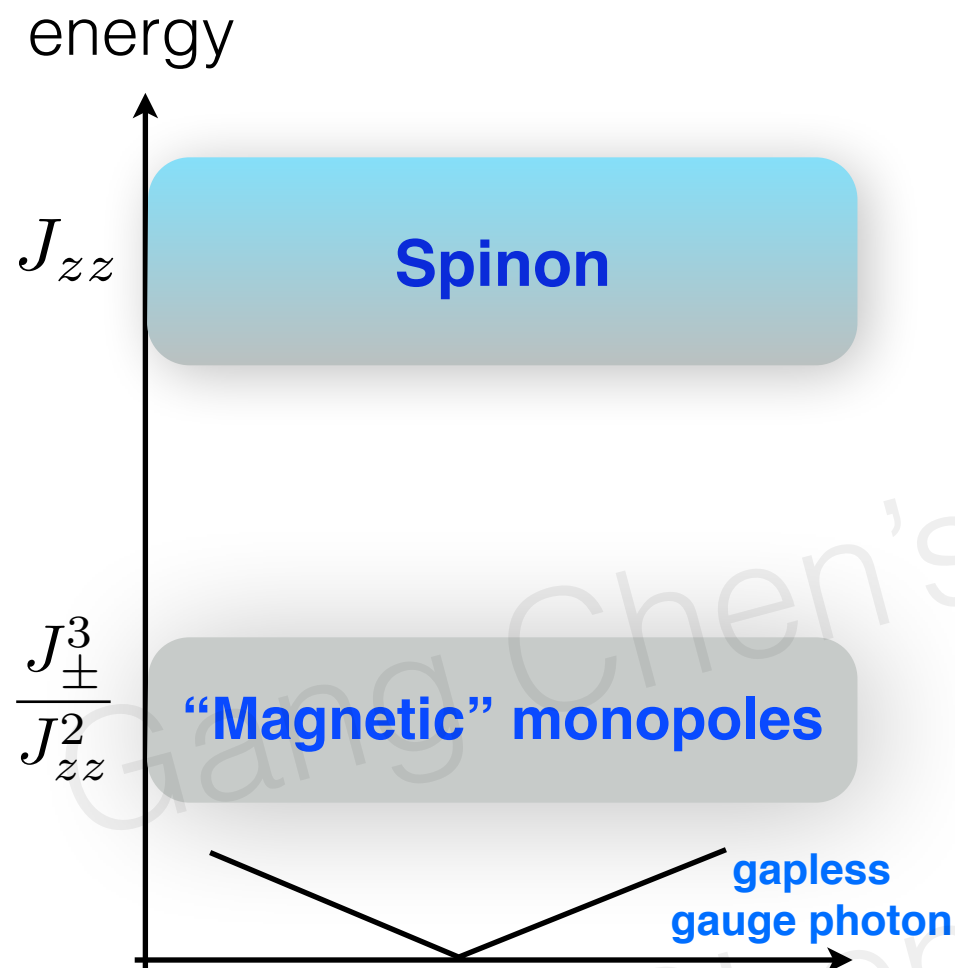
the right description of  
fractionalization and emergent

there are 3 elementary  
gauge photon,  
it is not goldstone boson  
breaking, there is no symmetry  
of emergent gauge theory

there are deconfined  
create 2 spinons,  
you can flip the further  
arbitrary distance, it is  
image there is a string  
spin is strongly fluctuating  
spinon are deconfined

the 3rd is magnetic monopole  
defect of the emergent

It is exotic phase of matter  
properties as topological order



Figs from Moessner

**Spinon deconfinement**

Emergent electric field

$$S^z \sim E$$

Emergent vector potential

$$S^\pm \sim e^{\pm iA}$$

$$H = J_{zz} \sum_{\langle i,j \rangle} S_i^z S_j^z - J_{\pm} \sum_{\langle i,j \rangle} (S_i^+ S_j^- + S_i^- S_j^+) + \dots$$

# Field-driven Higgs transition

How to tell if Ce2Sn2O7 is an octupolar U(1) QSL or not ?

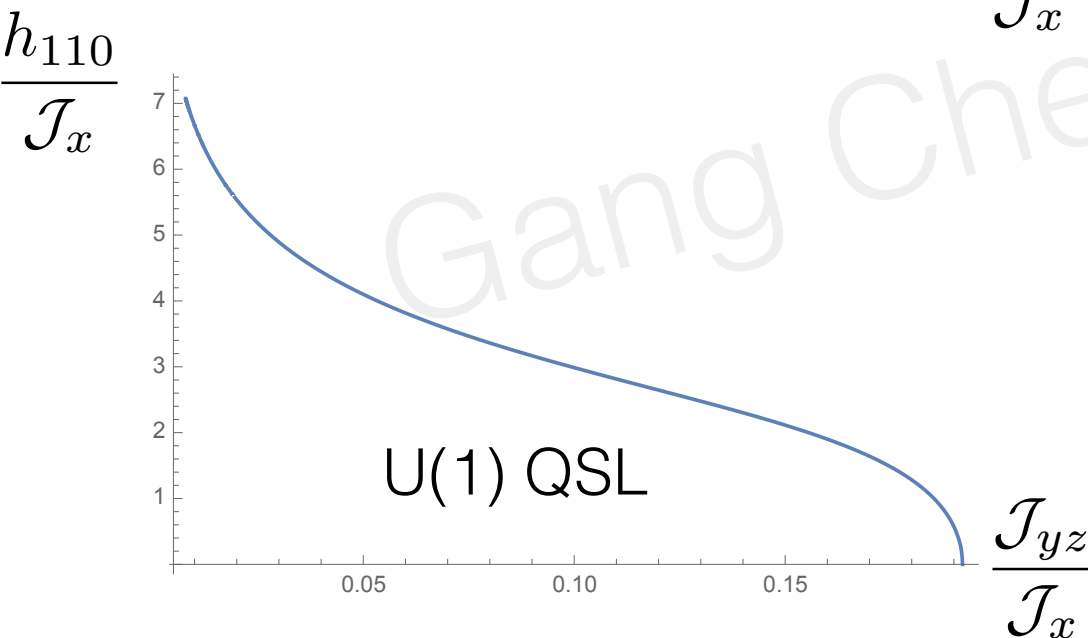
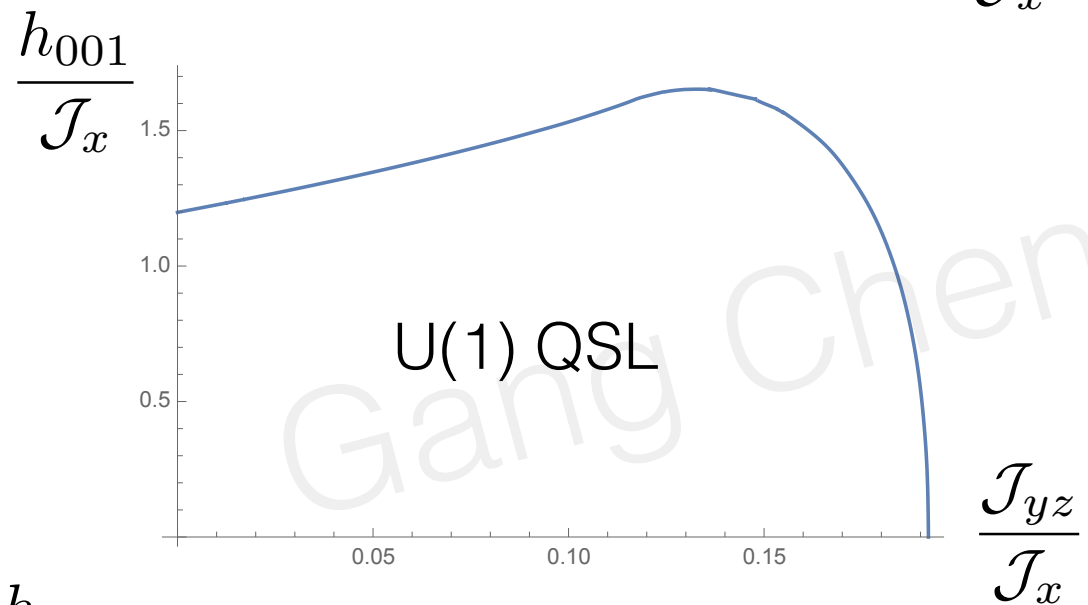
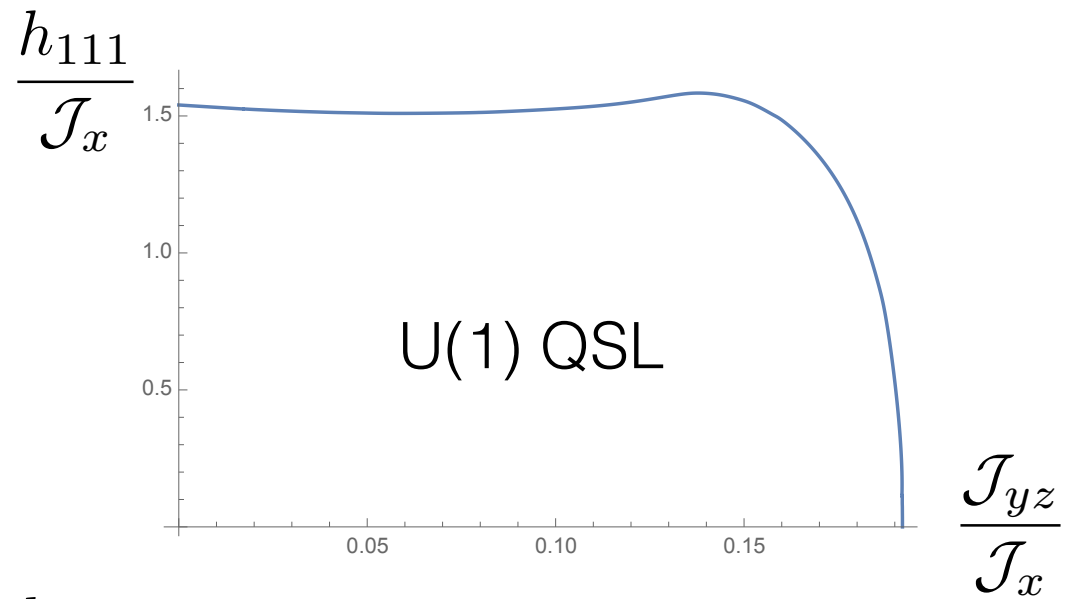
The idea to use a little knob that could simply lead to some clear experimental consequence, very much like the isotope effect of BCS supercond.

Here we apply external magnetic field, and expect a field-driven Higgs transition to magnetic ordering as the **field only couples to the matter field** (spinons).

$$H = \sum_{\langle ij \rangle} \mathcal{J}_x \tau_i^x \tau_j^x - \mathcal{J}_{yz} (\tau_i^y \tau_j^y + \tau_i^z \tau_j^z) - h \sum_i \tau_i^z (\hat{n}_i - 1)$$

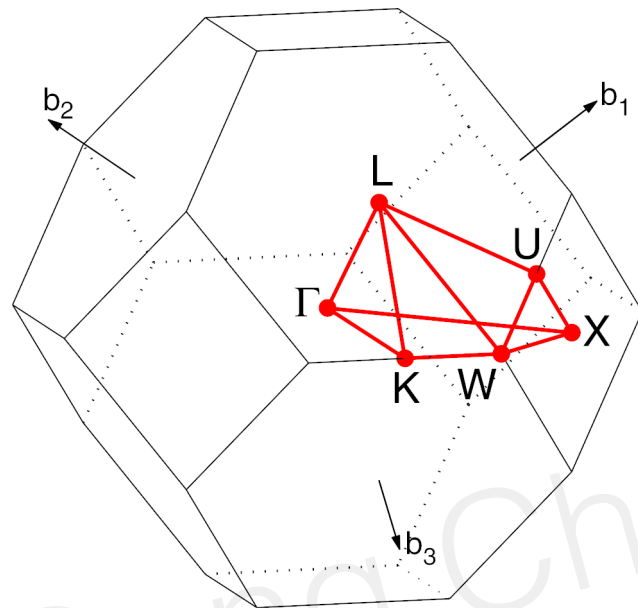
without losing generality

Higgs transition is very much like meissner effect in superconductor





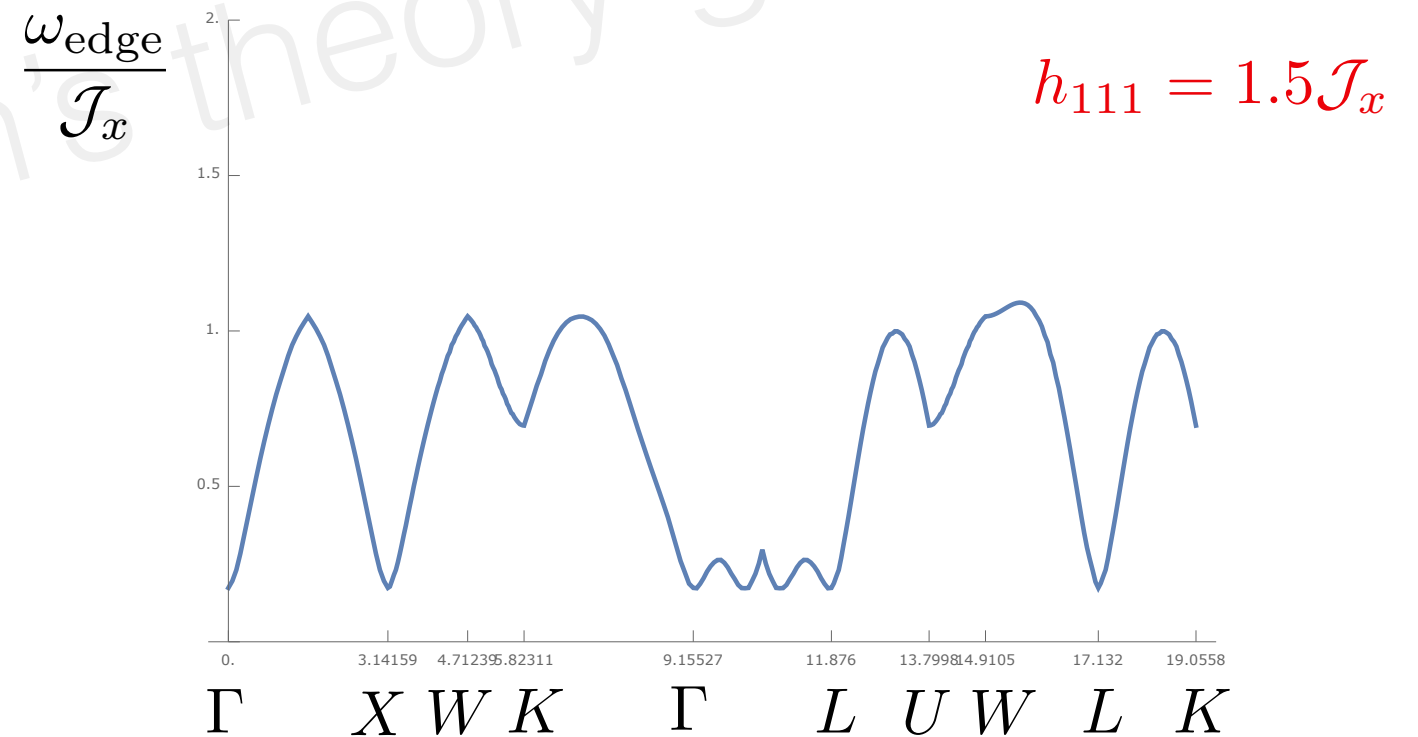
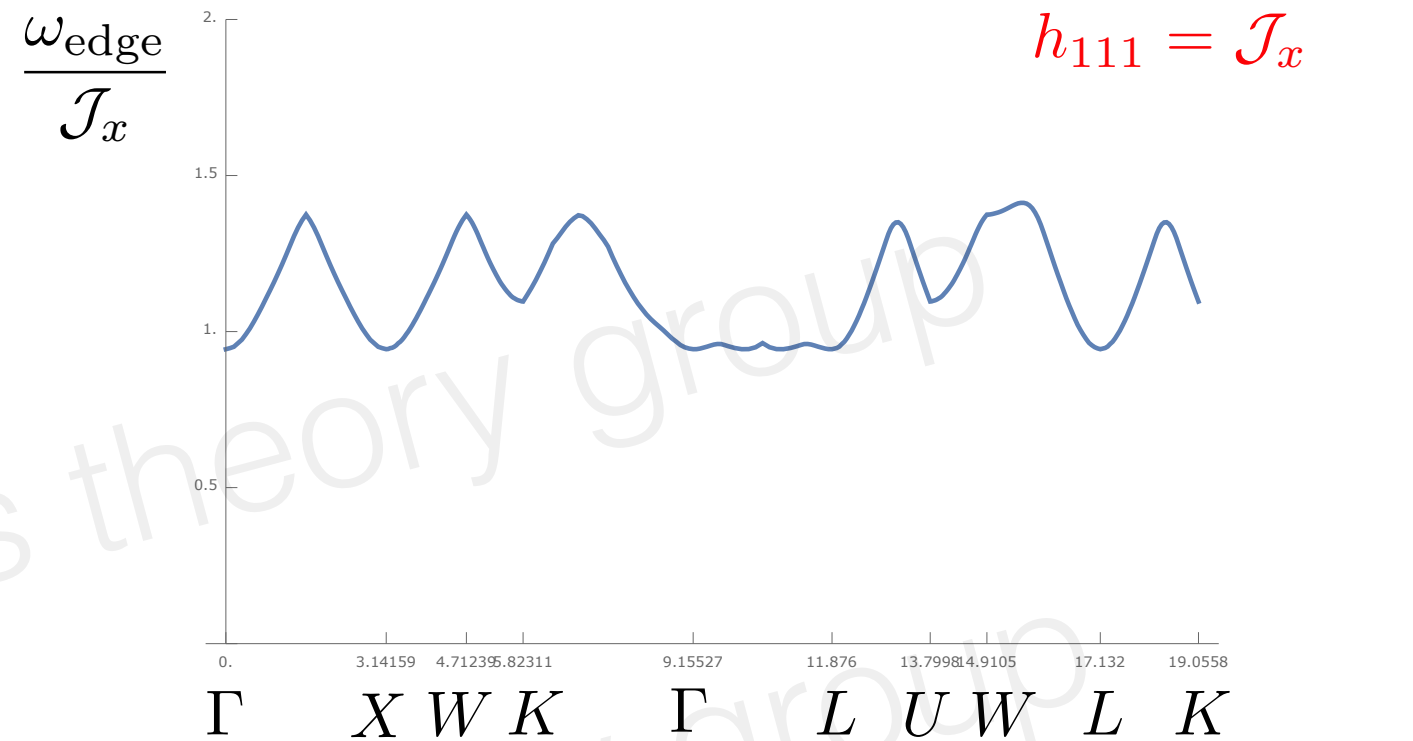
# Lower excitation edge



FCC path:  $\Gamma$ -X-W-K- $\Gamma$ -L-U-W-L-K|U-X

[Setyawan & Curtarolo, DOI: 10.1016/j.commatsci.2010.05.010]

Brioullin zone of  
FCC lattice



# Neutron scattering and thermal transport

## Dipolar- $U(1)$ QSL

neutron spin couples to both gauge field and matter field, observe both gapless gauge photon and gapped spinon continuum

## Octupolar- $U(1)$ QSL

neutron spin only couples to the matter field (spinons), observes only the gapped spinon continuum. External magnetic field can manipulate the spinon continuum, which can be confirmed by neutron scattering.

## Thermal transport

see both contribution, but there is a big separation of energy scales in spinon and gapless photons.

It can be beneficial to observe the low temperature peak in the thermal transport.

# Summary

We propose a general mechanism due to “weak” crystal electric field for the emergent magnetic order.

We propose a field-driven Anderson-Higgs’ transition as a simple knob to identify the octupolar  $U(1)$  quantum spin liquid.

Thank you !