

# Connect Emergence to Reality:

“Magnetic Monopole” Condensation out of  
 $U(1)$  Topological Order

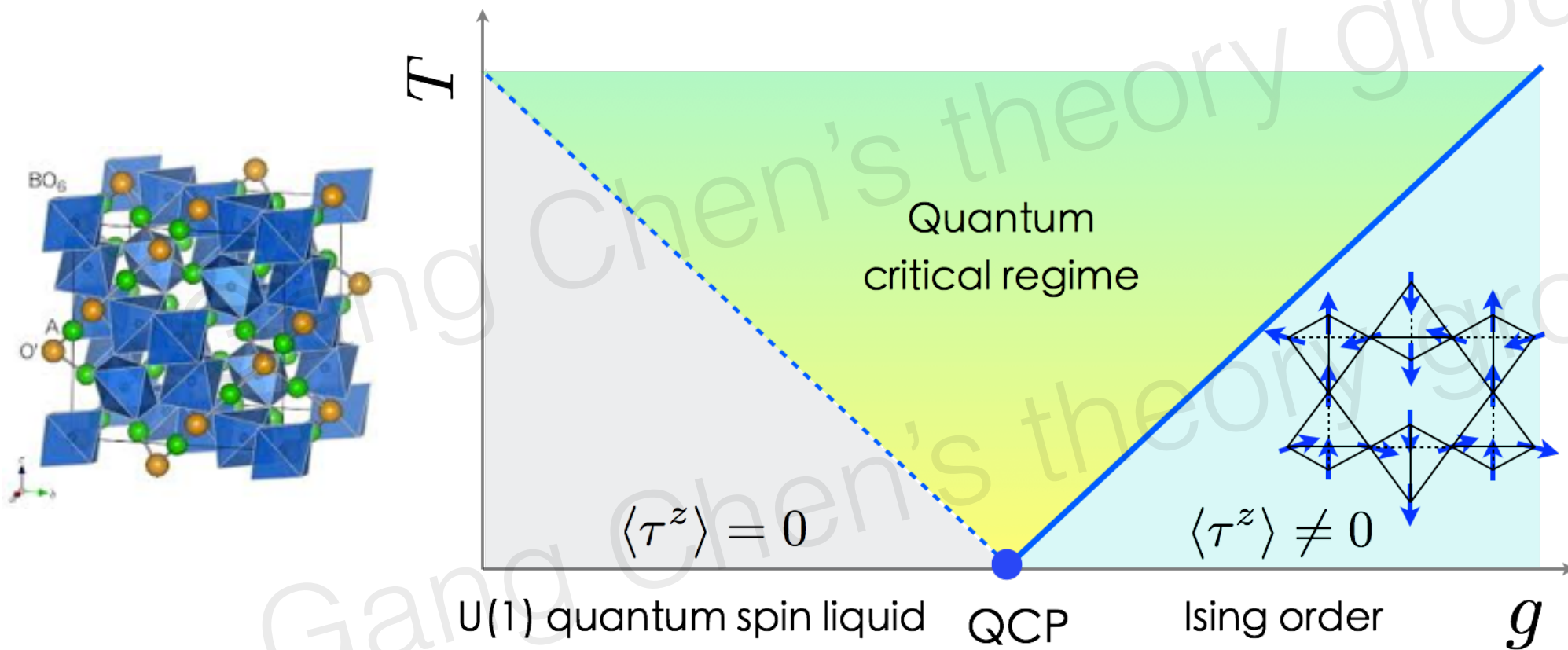
Gang Chen  
Fudan University



# OUTLINE

Monopole condensation transition out of U(1) topological order.

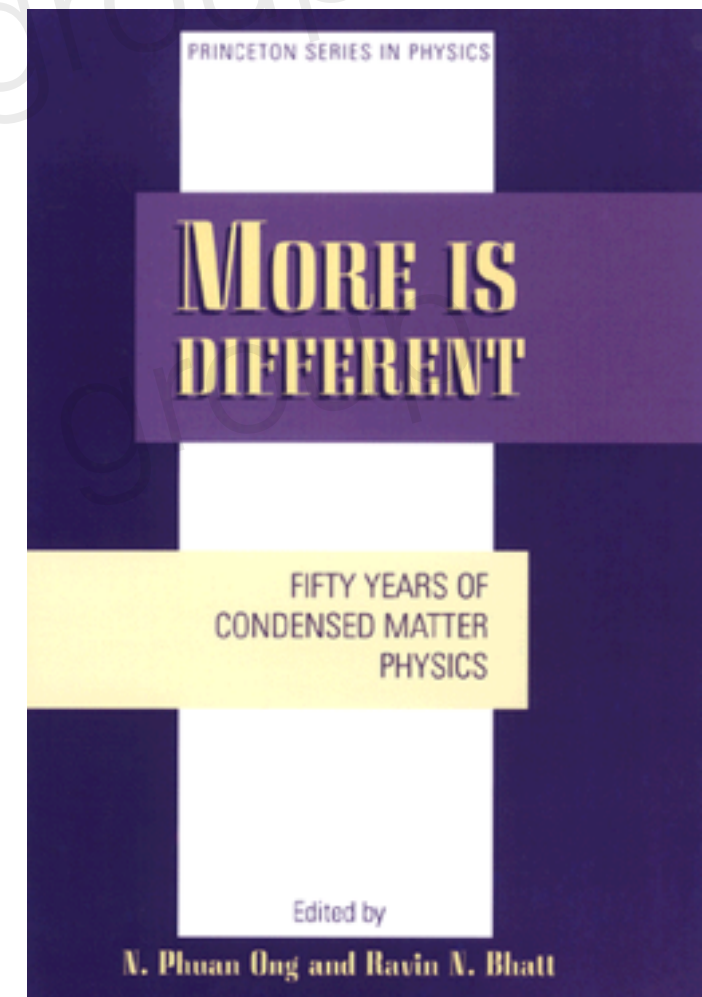
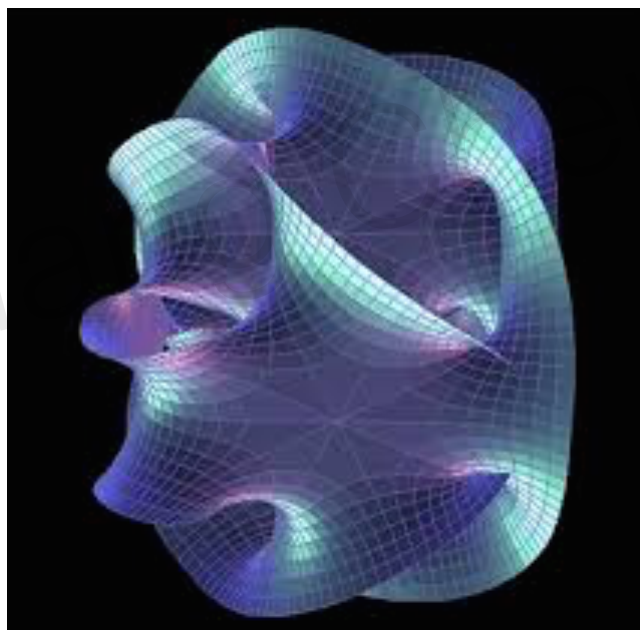
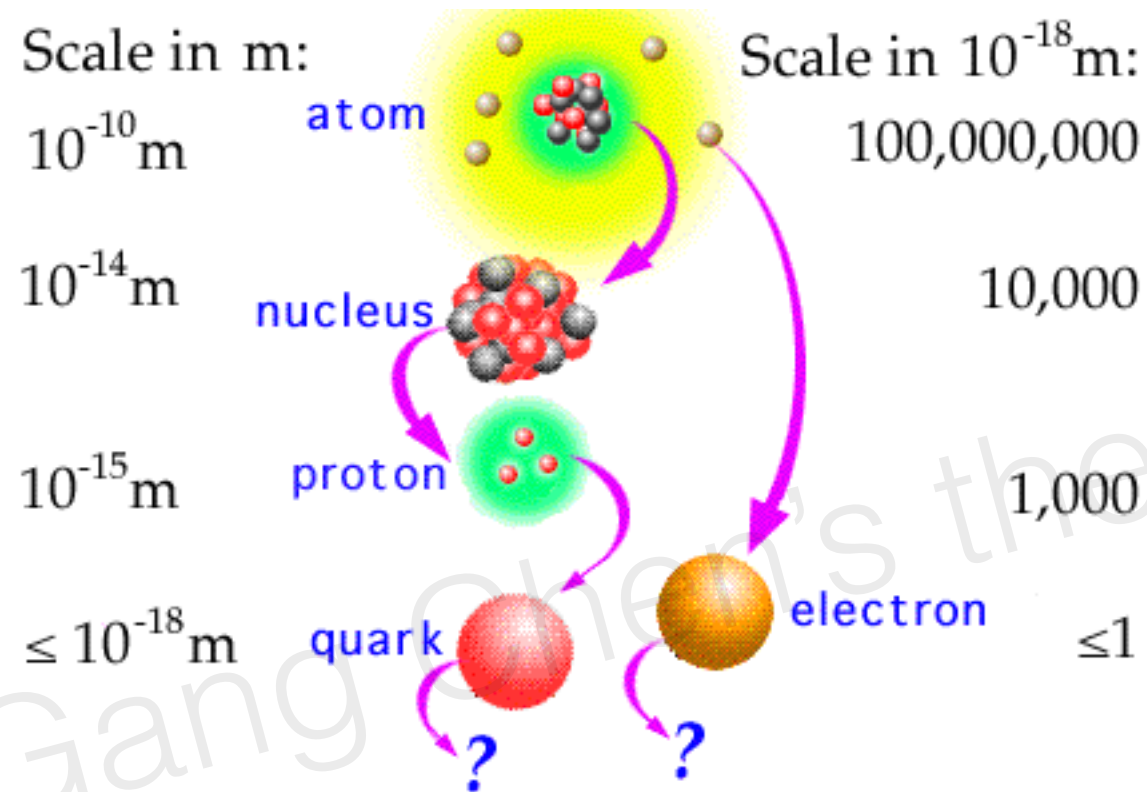
- I propose  $\text{Pr}_2\text{Ir}_2\text{O}_7$  sample is close to quantum phase transition between a 3D U(1) topological ordered state and Ising order.



**Gang Chen**, arXiv 1602.02230, Phys. Rev. B. **94**, 205107 (2016)

**Yao-Dong Li, Gang Chen**, in preparation, 2017  
for conduction electrons.

# Reduction vs Emergence

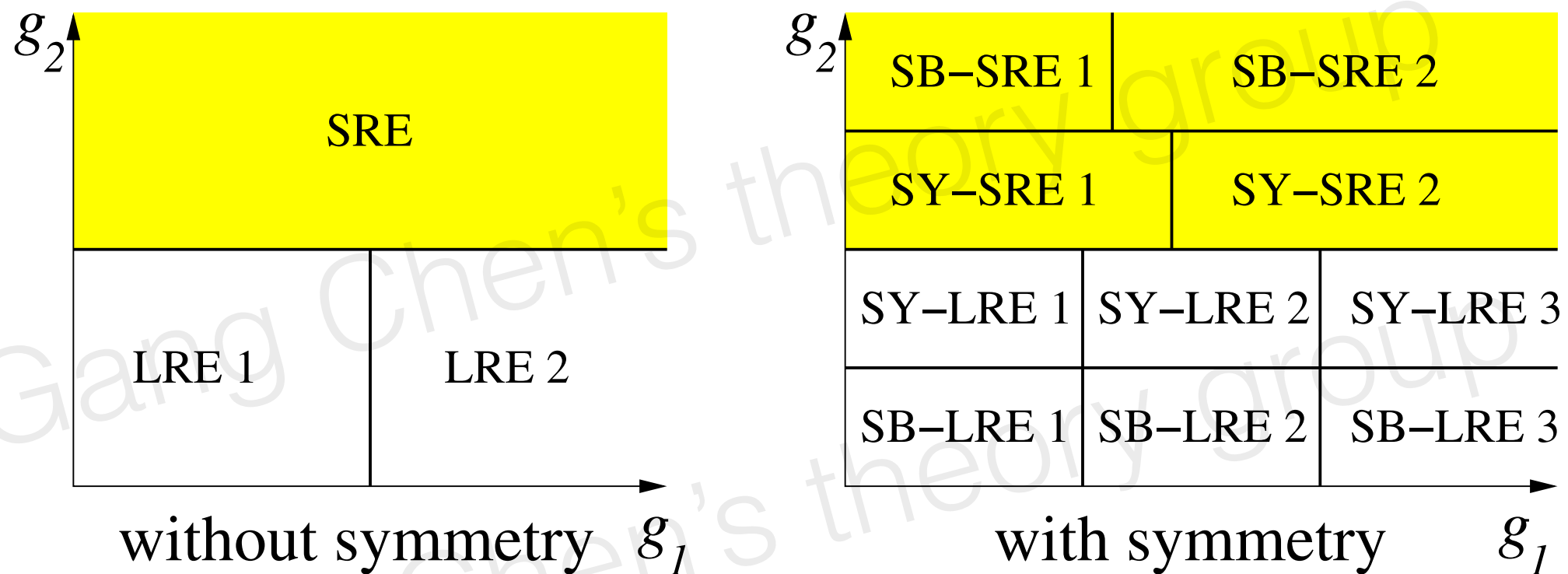


Condensed matter is full of emergence



# Wen's universal phase diagram for symmetry and topology

Xiao-Gang Wen



LRE=long-range entanglement  
SRE=short-range entanglement

**Intrinsic topological order has long range entanglement (LRE).**

the interplay between  
great discovery of  
symmetry breaking  
entanglement.

For intrinsic topological  
topological order, the

symmetry and intrinsic

finite step local unitary

The classification of  
the degree of freedom  
understanding how

Here I will give a review

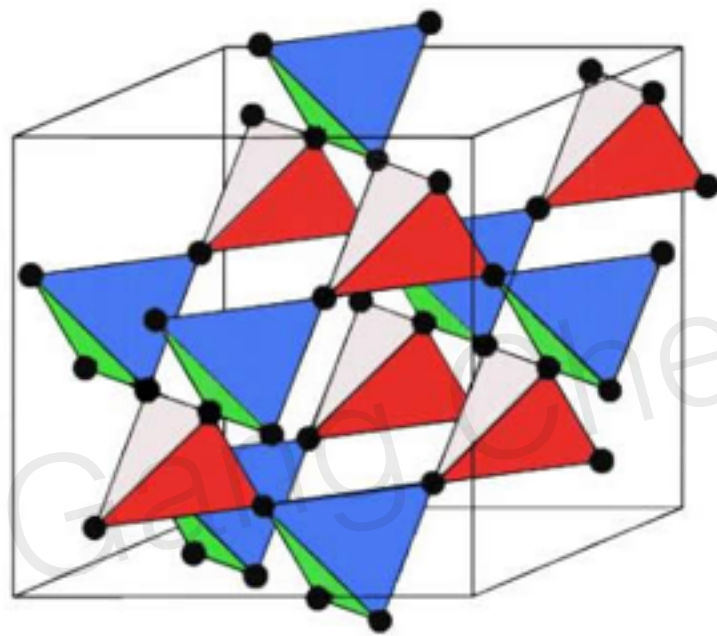
finite step local unitary



# Monopole condensation out of $U(1)$ topological order

- Introduction to spin ice, classical and quantum.
- Magnetic transition of quantum spin ice  $U(1)$  quantum spin liquid is the confinement transition of compact  $U(1)$  lattice gauge theory (or compact quantum electrodynamics)
- Monopole condensation and proximate phases

# Spin ice in rare-earth pyrochlores



## Rare Earth Elements

by Geology.com

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Over years, there are a lot of activity in spin ice system.

spin ice is realized in rare earth pyrochlore systems, where the rare earth ions form pyrochlore lattice and host the Ising spin. because of the crystal field effect, the Ising spin points either into or out of the center of the tetrahedron

The interaction between the ising is AFM, it favor 2 spin in 2 spin out of the tetrahedra. This is the 2-in 2-out spin ice rule.

Beucase of the analog relation with H position in water ice, each O has 4 H near it, 2 are close, 2 are further.

# Spin ice in rare-earth pyrochlore

$$H = J_{zz} \sum_{\langle i,j \rangle} S_i^z S_j^z + \text{dipolar}$$

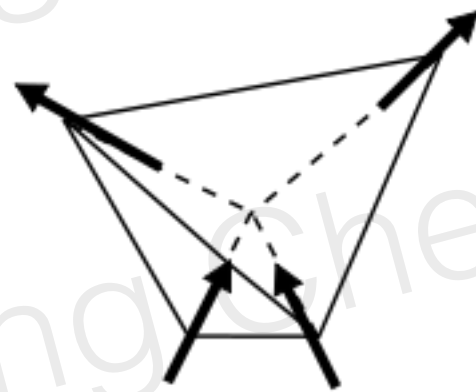
Castelnovo, Gingras, Moessner, Sondhi, Schiff  
Penc, .....

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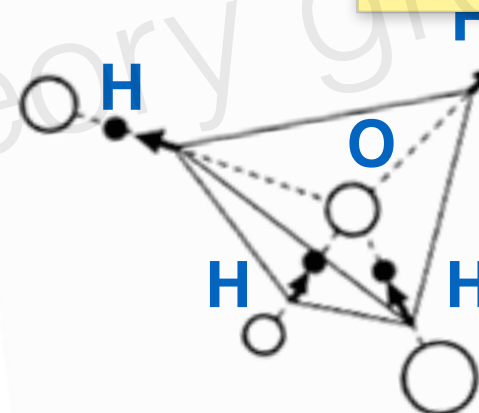
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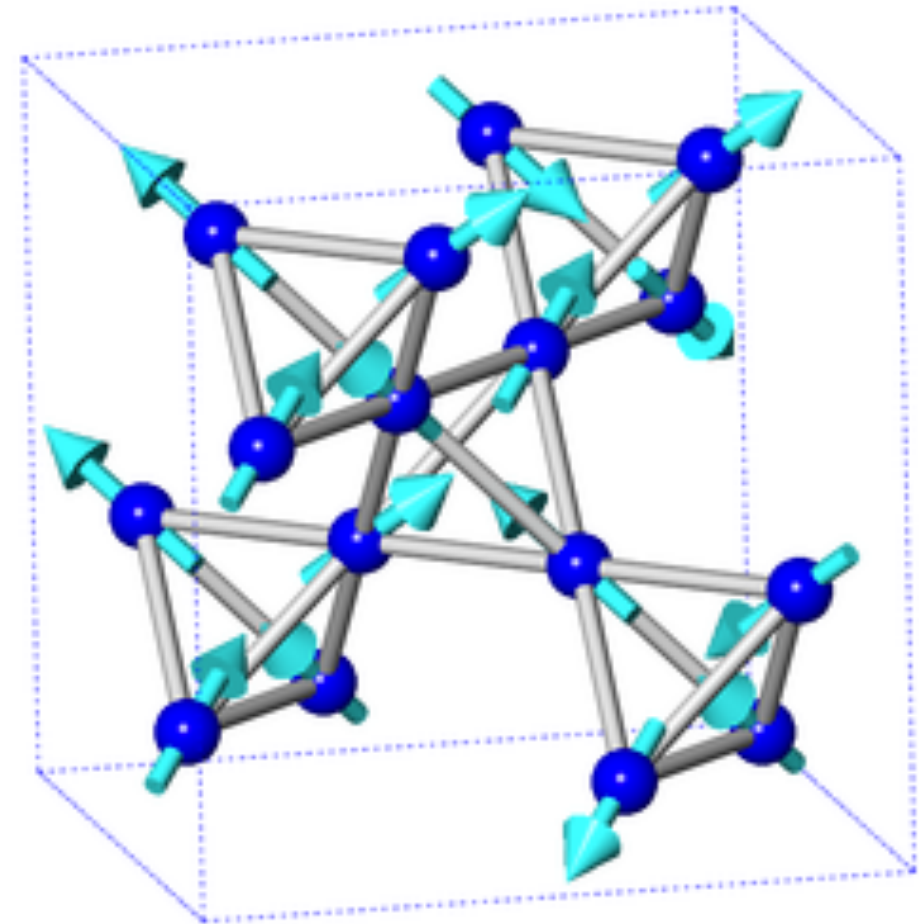
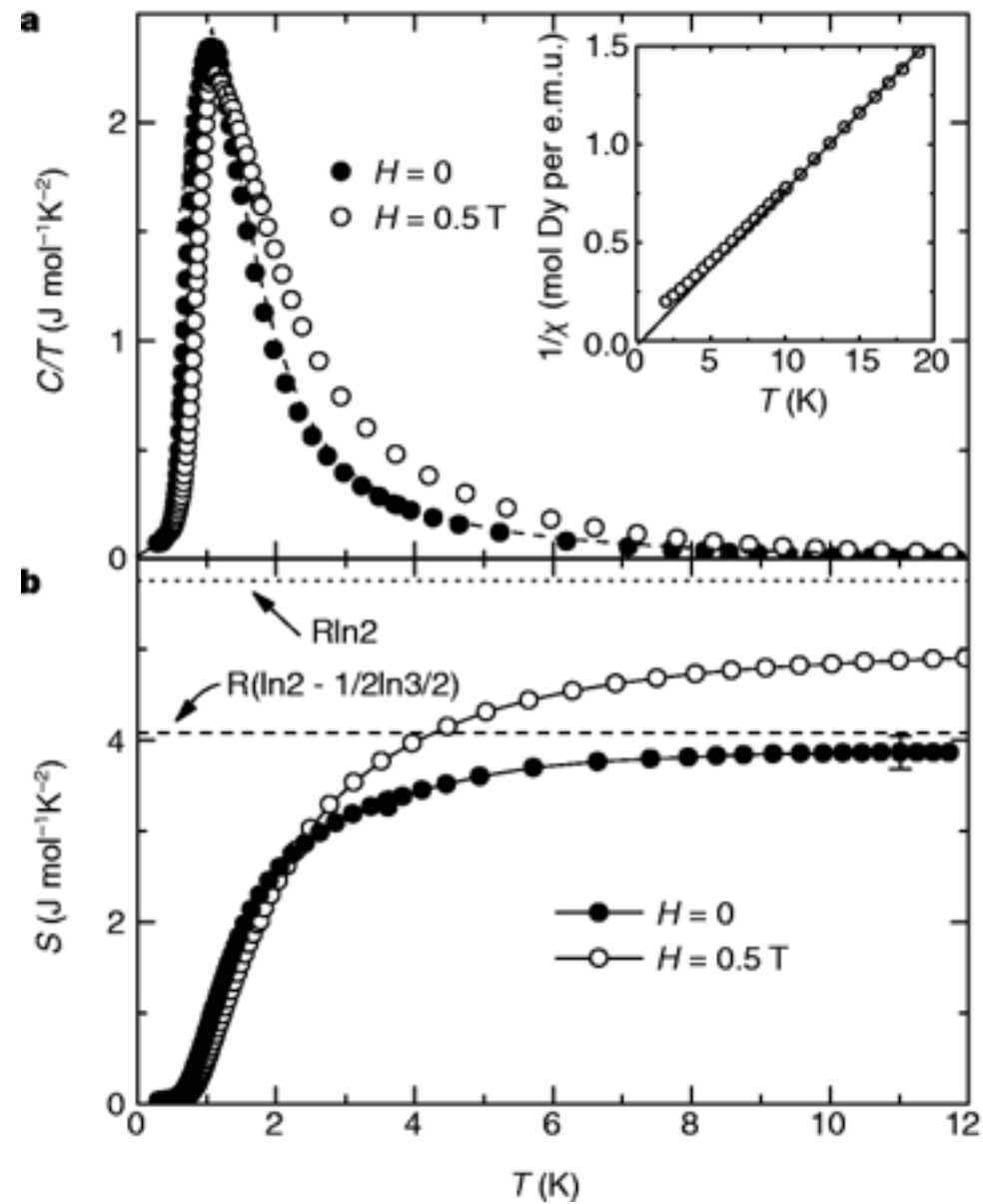
2-in 2-out  
**spin ice rule**



2-in 2-out  
**water ice rule**

from wiki

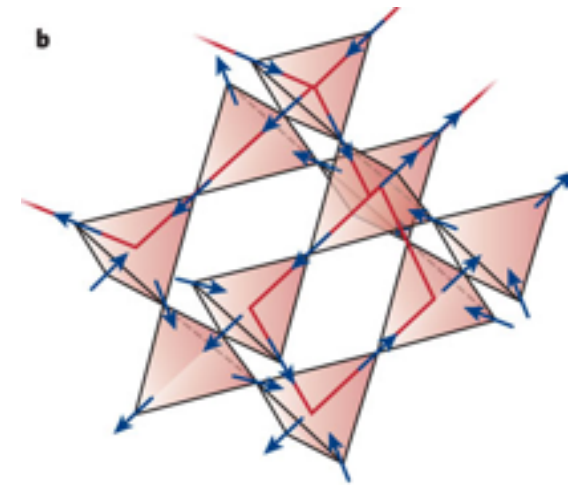
# Classical spin ice



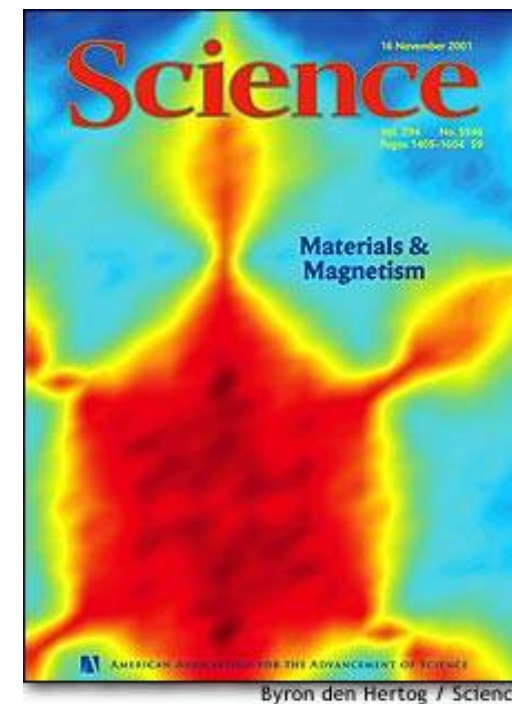
Pauling entropy in spin ice,  
Ramirez, etc, Science 1999



# Classical spin ice



- The “2-in 2-out” states are extensively degenerate.
- At  $T < J_{zz}$ , the system **thermally** fluctuates within the ice manifold, leading to classical spin ice and interesting experimental discoveries.



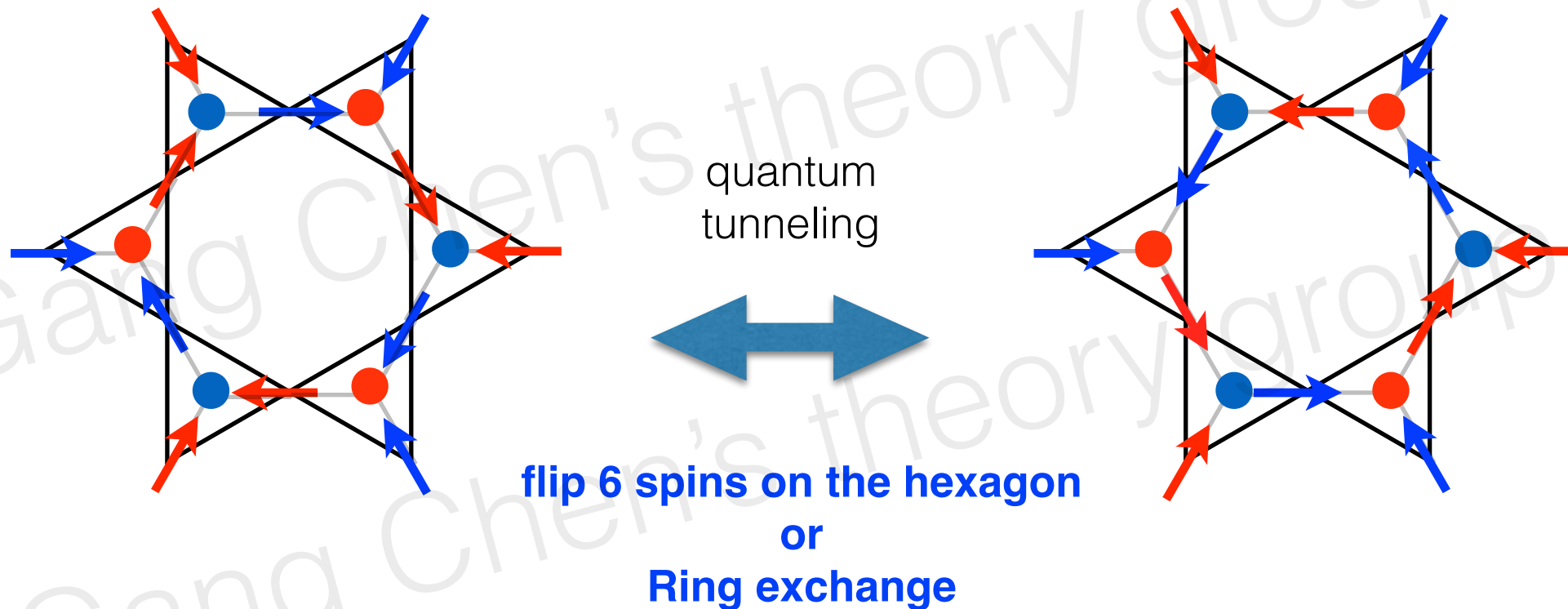
1. the 2-in 2-out spin ice is actually extensively degenerate. On each tetrahedron, one can choose either the 2-in or the other 2-out.
2. at  $T \ll J_{zz}$ , the system is thermally fluctuating within the ice manifold, leading to classical spin ice and interesting experimental consequences. Many of them have been observed in nature and science.

Pinch points in spin correlation

# Quantum fluctuation can lead to U(1) QSL

$$H = J_{zz} \sum_{\langle i,j \rangle} S_i^z S_j^z - J_{\pm} \sum_{\langle i,j \rangle} (S_i^+ S_j^- + S_i^- S_j^+) + \dots$$

Hermele, Fisher, Balents,  
Moessner, Isakov, YB Kim....

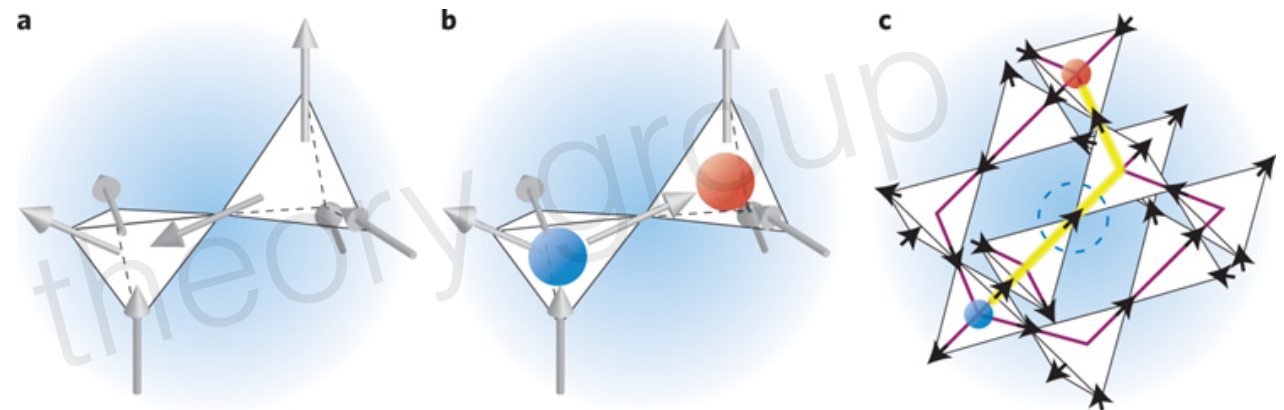
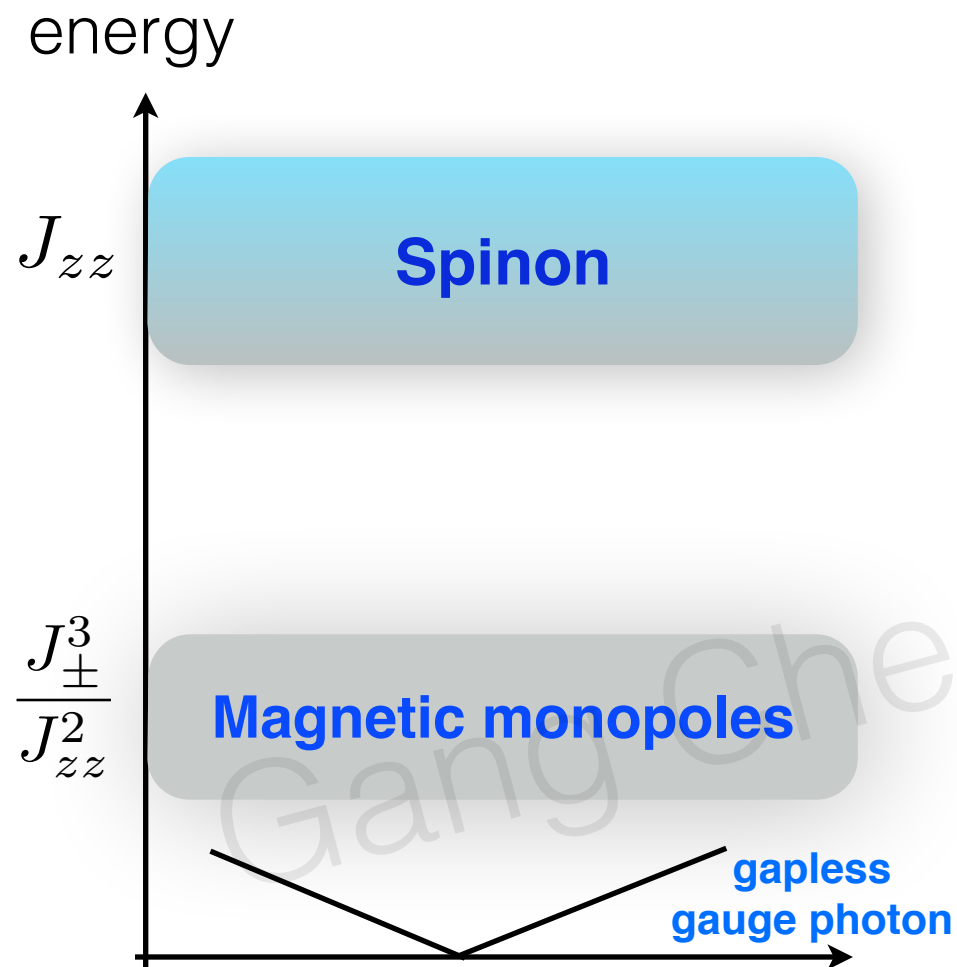


- Pretty much one can add any term to create **quantum** tunneling, as long as it is not too large to induce magnetic order, the **ground state** is a quantum spin ice !

1. But classical spin ice is purely classical and is not a new phase of matter. It is smoothly connected to high temperature paramagnetic phase.

2. In contrast, quantum spin ice is a new quantum phase of matter.

# U(1) QSL is NOT a Landau symmetry breaking phase



Figs from Moessner&Schiffer,2009

## Spinon deconfinement

- Unlike CSI, QSI is a **novel phase of matter**. No LRO, no symmetry breaking, **cannot** be understood in Landau's paradigm!
- The right description is in terms of fractionalization and emergent gauge structure.

as quantum spin ice is a disordered state, there is no long range order, no symmetry breaking, it is a new phase of matter and cannot be understood in the Landau's paradigm of symmetry breaking.

Important question: Has 3D  $U(1)$  QSL been realized in experiments,  
or realized in the context of spin ice?

What would be the experimental evidence?

one may wonder if qsi exist in some physical system.

The answer is probably.

one can write a realistic hamiltonain and show, (even prove) the ground state should be quantum spin ice.

the real difficulty is to confirm it experimentally.

because it does not have LRO, unlike trivial order phase, it is very difficult to confirm it.

TO cofnirm it ,one should observe either deconfined spinons or emergent gless



# Realistic models

- Kramers' doublet

$$H = \sum_{\langle ij \rangle} \{ J_{zz} \mathbf{S}_i^z \mathbf{S}_j^z - J_{\pm} (\mathbf{S}_i^+ \mathbf{S}_j^- + \mathbf{S}_i^- \mathbf{S}_j^+) \\ + J_{\pm\pm} (\gamma_{ij} \mathbf{S}_i^+ \mathbf{S}_j^+ + \gamma_{ij}^* \mathbf{S}_i^- \mathbf{S}_j^-) \\ + J_{z\pm} [\mathbf{S}_i^z (\zeta_{ij} \mathbf{S}_j^+ + \zeta_{ij}^* \mathbf{S}_j^-) + i \leftrightarrow j] \},$$

S. H. Curnoe, PRB (2008).

Savary, Balents, PRL 2012

- Non-Kramers' doublet

$$H = \sum_{\langle ij \rangle} \{ J_{zz} \mathbf{S}_i^z \mathbf{S}_j^z - J_{\pm} (\mathbf{S}_i^+ \mathbf{S}_j^- + \mathbf{S}_i^- \mathbf{S}_j^+) \\ + J_{\pm\pm} (\gamma_{ij} \mathbf{S}_i^+ \mathbf{S}_j^+ + \gamma_{ij}^* \mathbf{S}_i^- \mathbf{S}_j^-) \}$$

S. Onoda, etc, 2009

SB Lee, Onoda, Balents, 2012

- Dipole-octupole doublet

$$H = \sum_{\langle ij \rangle} J_x S_i^x S_j^x + J_y S_i^y S_j^y + J_z S_i^z S_j^z \\ + J_{xz} (S_i^x S_j^z + S_i^z S_j^x).$$

Y-P Huang, **Gang Chen**, M Hermele, PRL 2014

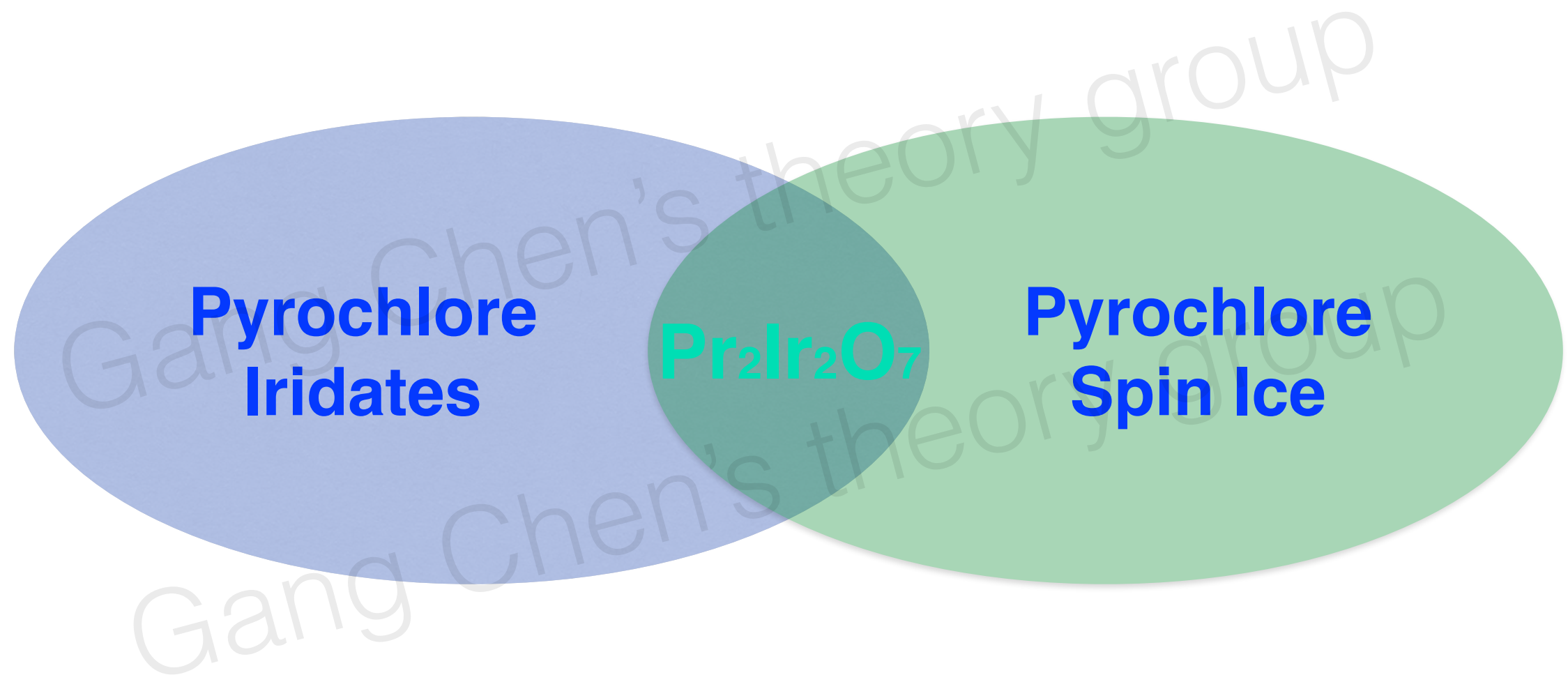
Yaodong Li, **Gang Chen**, Arxiv 1607

Nd<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub>, Nd<sub>2</sub>Sn<sub>2</sub>O<sub>7</sub>, Nd<sub>2</sub>Zr<sub>2</sub>O<sub>7</sub>, Ce<sub>2</sub>Sn<sub>2</sub>O<sub>7</sub>, etc

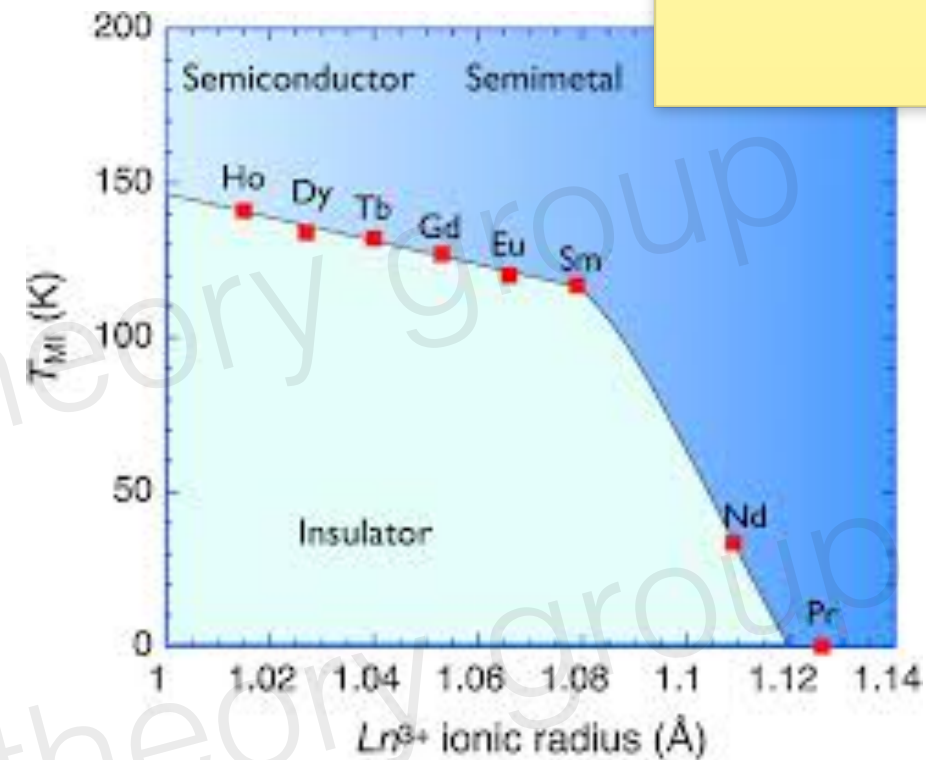
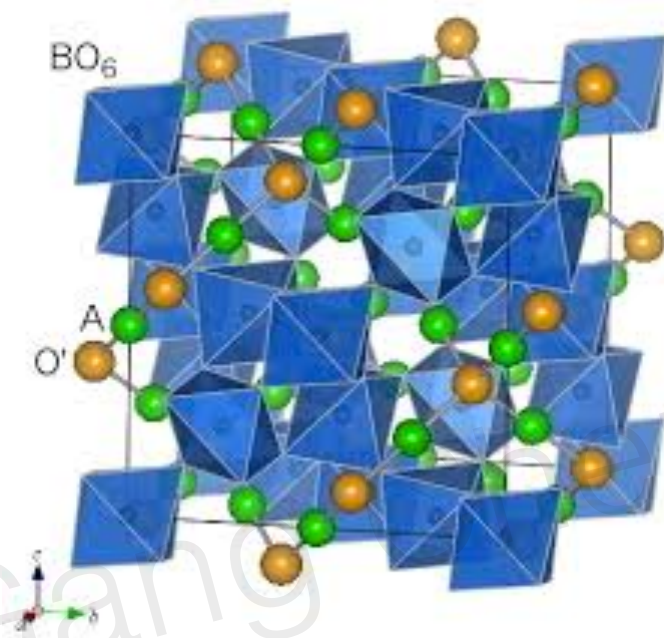
no sign problem for QMC on any lattice.

It supports nontrivial phase like quantum spin ice U(1) quantum spin liquid.

# Pyrochlore Iridate and Pyrochlore Spin Ice



# Pyrochlore iridates: $\text{Pr}_2\text{Ir}_2\text{O}_7$



K Matsuhira, M Wakeshima, Y Hinatsu, S. Takagi  
JPSJ, 2011

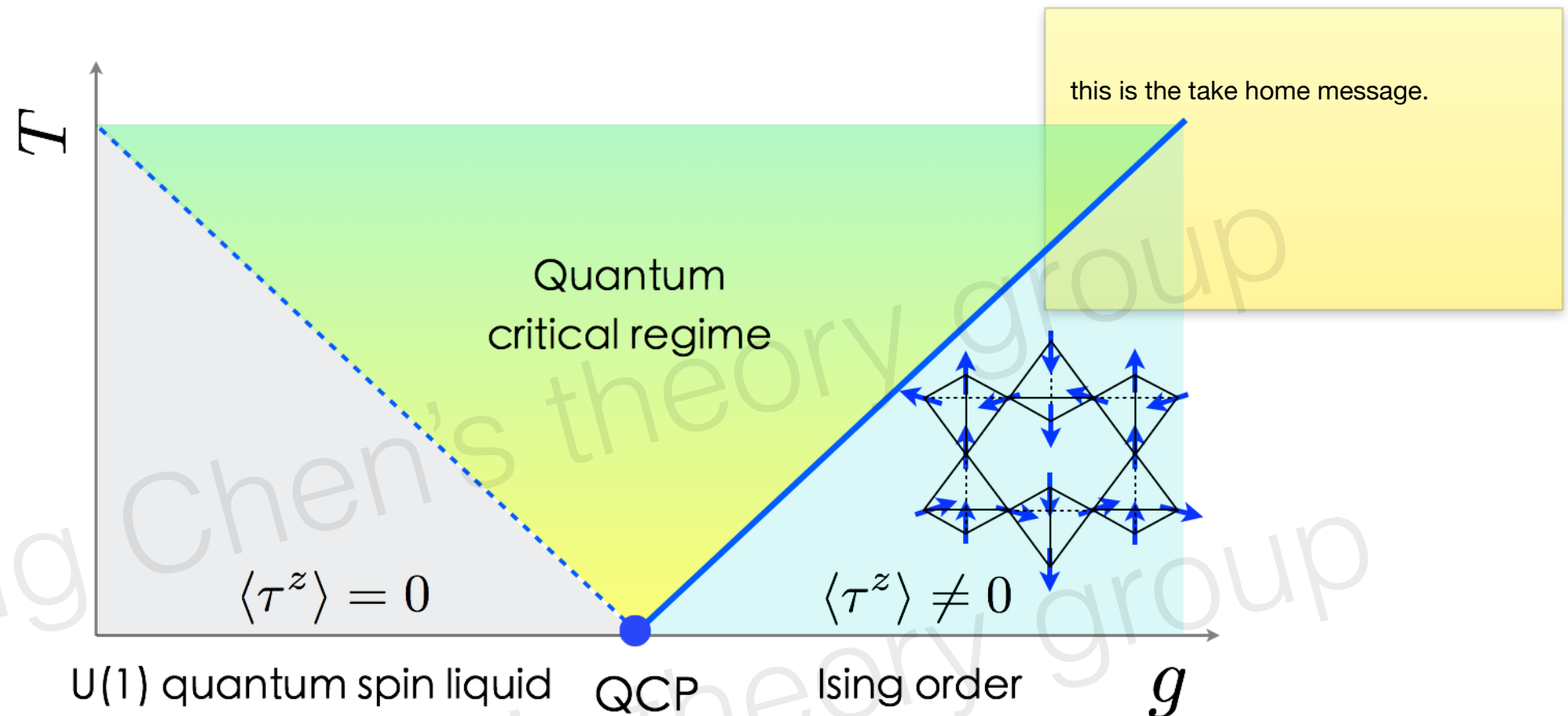
Many nice experimental works by S Nakatsuji, P Gegenwart, L Balicas, etc

Ref: D Pesin, L Balents, 2009, Xian-Gang Wan, etc 2010, Witczak-Krempa, Yong Baek Kim, SungBin Lee;  
Michael Hermele, Gang Chen, etc

Over the years, there are a lot of effort in pyrochlore iridates, almost all of them are insulator with some magnetic order, except

$\text{Pr}_2\text{Ir}_2\text{O}_7$  is unique, it remains metallic, and so most of the work in the field focus on the iridates to discuss local moment physics.

# My proposal for $\text{Pr}_2\text{Ir}_2\text{O}_{7-\delta}$

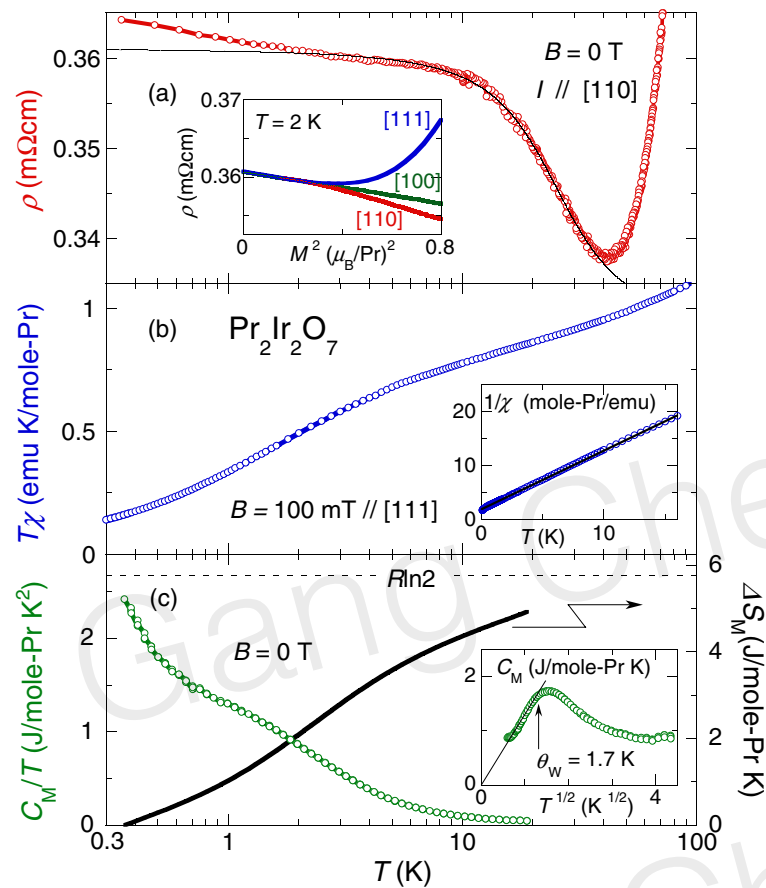


Pr local moments are close to a “magnetic” monopole condensation transition from quantum spin ice quantum spin liquid to an AFM long-range ordered state.

The Ir conduction electrons may drive the transition, but do not influence the nature of the phase transition.

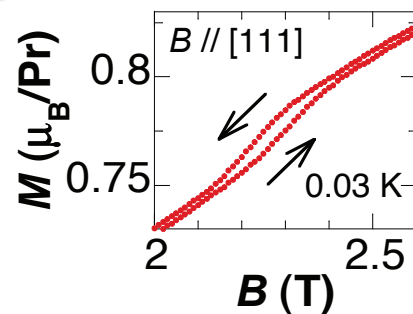


# Pr<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub> a featureless disordered state near an ordered state



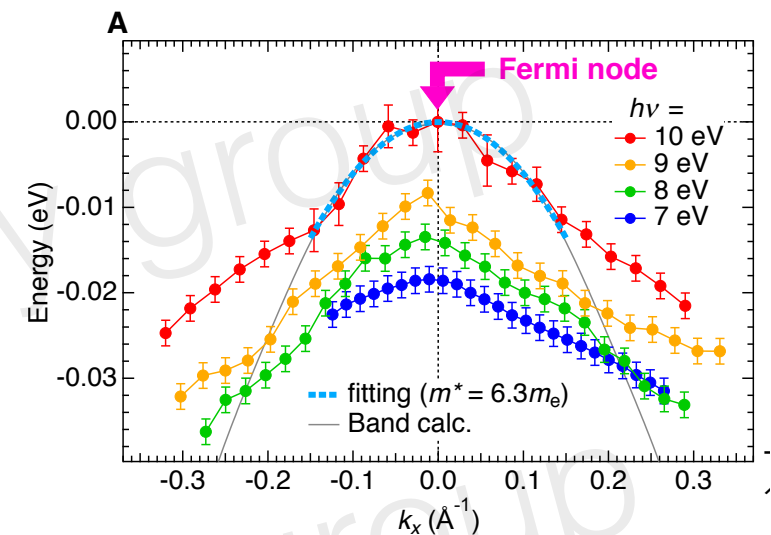
Nakatsuji, etc

PRL **96**, 087204 (2006)



metamagnetic transition

ARPES: quadratic band touching of Ir 5d electrons



Valence band approaches the Fermi energy at few meV resolution

T Kondo, S Shin, etc 2014

B J Yang, Yong Baek Kim 2011

E G Moon, CK Xu, Y B Kim, L Balents, 2013

Expts are **sample dependent**,  
Recently, some samples are found AFM ordered.

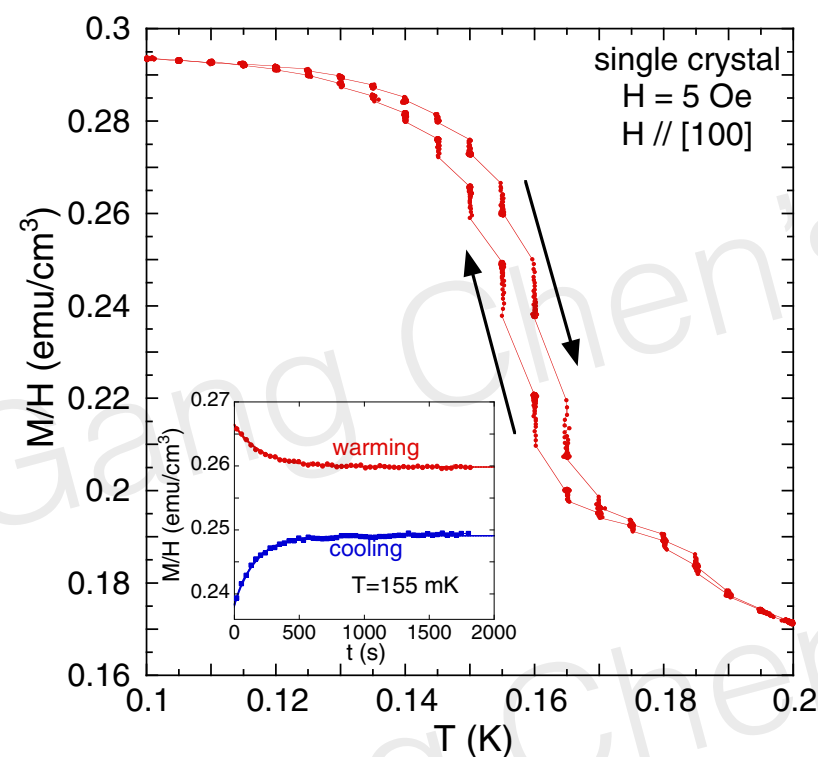
Yaodong Li, Gang Chen  
in **preparation 2017**

# Experiments: a featureless state near an ordered state

PHYSICAL REVIEW B **89**, 224419 (2014)

## First-order magnetic transition in $\text{Yb}_2\text{Ti}_2\text{O}_7$

E. Lhotel,<sup>1,\*</sup> S. R. Giblin,<sup>2</sup> M. R. Lees,<sup>3</sup> G. Balakrishnan,<sup>3</sup> L. J. Chang,<sup>4</sup> and Y. Yasui<sup>5</sup>



Some samples have FM LRO  
with 1st transition. Some samples  
do not have order.

PHYSICAL REVIEW X **1**, 021002 (2011)

## Quantum Excitations in Quantum Spin Ice

Kate A. Ross,<sup>1</sup> Lucile Savary,<sup>2</sup> Bruce D. Gaulin,<sup>1,3,4</sup> and Leon Balents<sup>5,\*</sup>

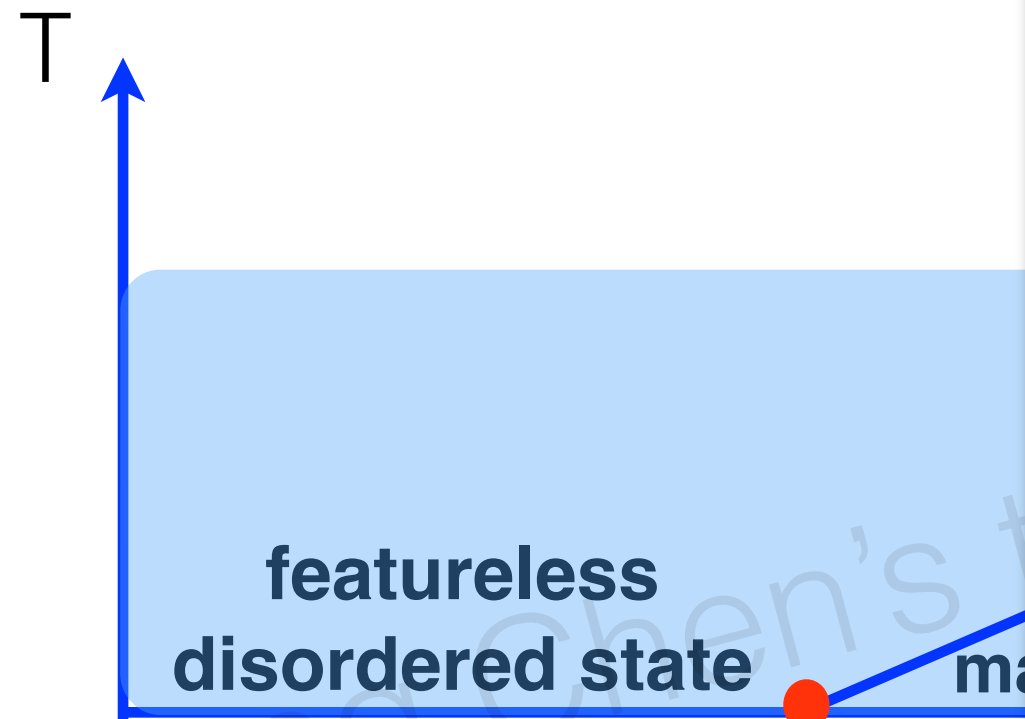
waves [14]. Although one neutron study [15] supported ferromagnetic order in  $\text{Yb}_2\text{Ti}_2\text{O}_7$ , intriguingly, the majority of neutron scattering measurements have reported a lack of magnetic ordering and the absence of spin waves at low fields in this material [16–18]. In a recent study,

this slide should be quick.

another system is purely local m  
also some sample order  
some sample do not order

but order ferromagnetically .

# Summary of experimental results



this is a summary of the messy experiments.

the system is probably near some transition between a featureless disordered state and magnetic ordered state.

and the tuning parameter can be chemical pressure, oxygen content.

important questions are

what is the phase transition?  
what is the nature of the phase transition?

disordered state is hard to characterize.

phase transition is more visible in experiments,

magnetic order is easy to detect in experiments.

One can use proximate phase and phase transition

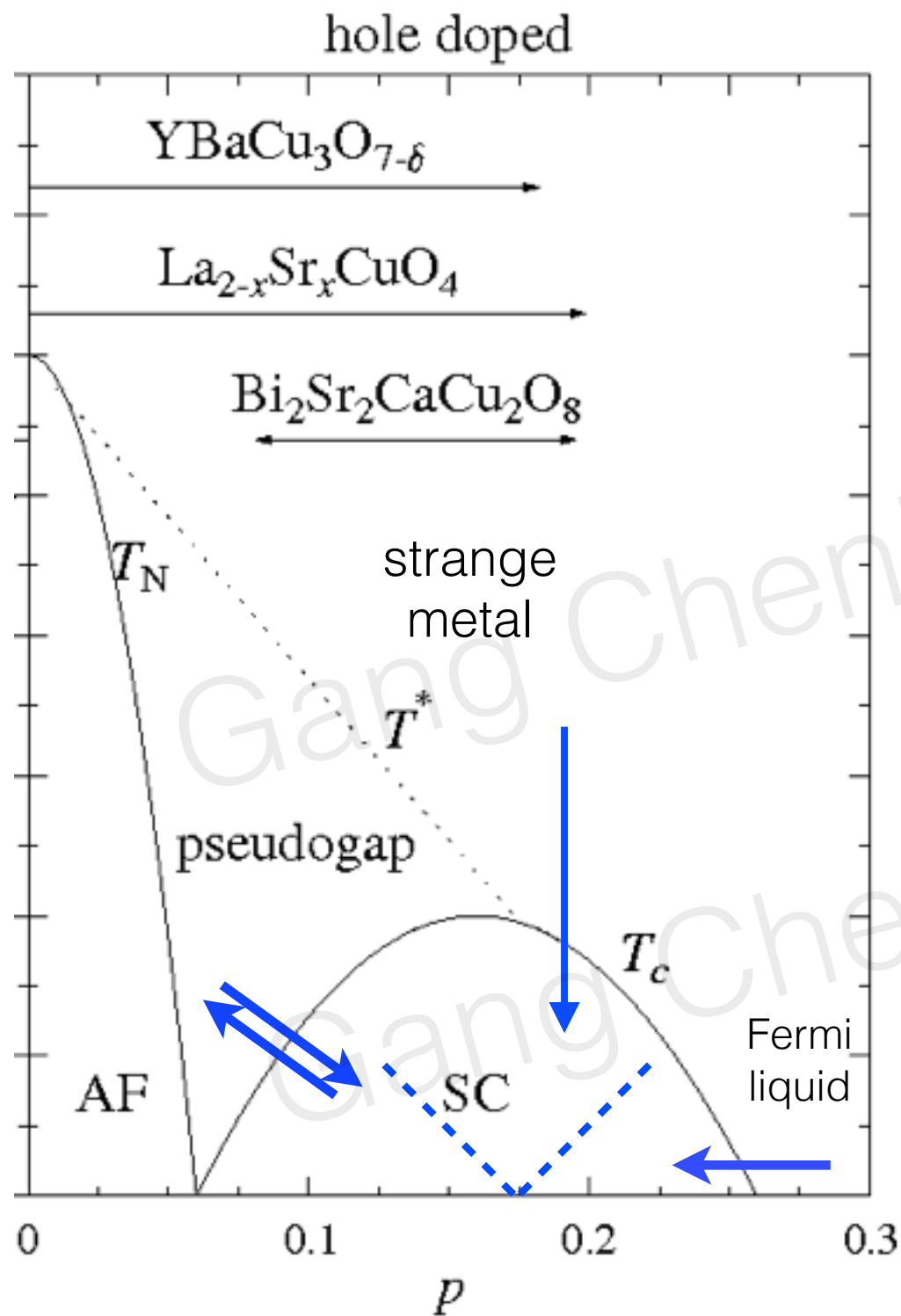
chemical pressure,  
oxygen content .....

- What is the structure of the magnetic order?
- What is the relationship between the disordered state and various magnetic states?
- What is the nature of the featureless disordered state?

ordered

Is it **QSI**?

# Insight from high-Tc superconductors



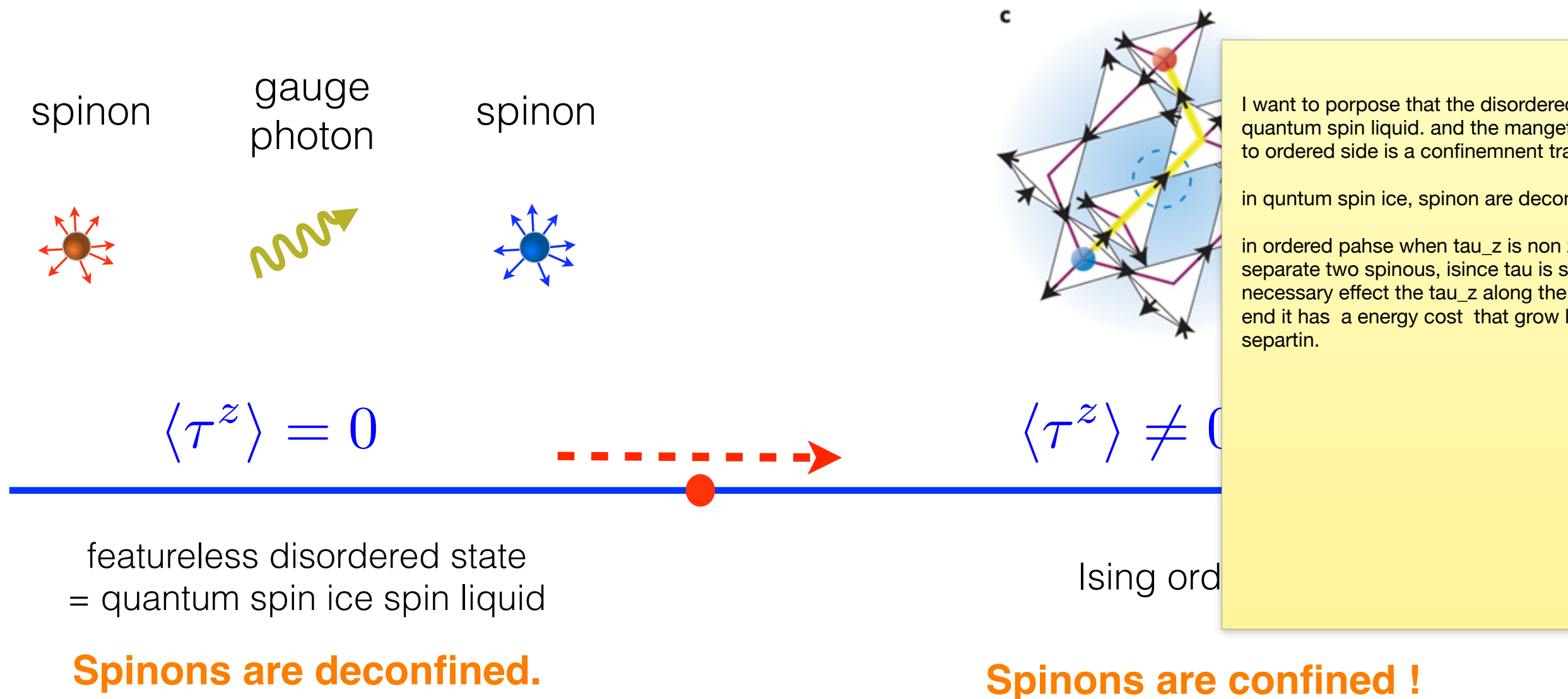
One important question is to understand the relationship between different phases (and/or orders)

1. Perturbative treatment (not interesting): instability of Fermi liquid;
2. Attack from top: instability of non-Fermi liquid;
3. Attack from Left, attack from Right: what is PG (Z2 topological order?) ? (Senthil, Balents, Nayak, Fisher 2000-2002);
4. Attack from bottom: some quantum criticality under the SC dome?

Figure from wiki

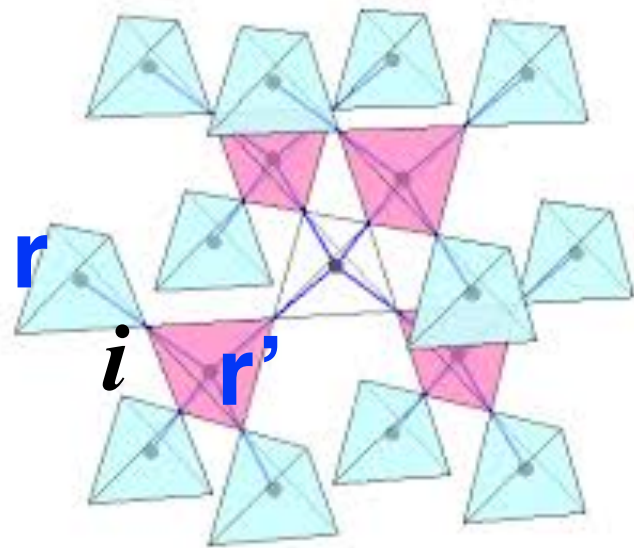


# Confinement transition out of U(1) quantum spin liquid



More generally, for **non-Kramers' doublet**, the magnetic transition out of U(1) QSL **MUST** be a confinement transition, this may apply to Tb<sub>2</sub>Ti<sub>2</sub>O<sub>7</sub>.

# Lattice gauge theory formalism: technique



$$E_{\mathbf{r}\mathbf{r}'} \sim \tau_i^z, e^{iA_{\mathbf{r}\mathbf{r}'}} \sim \tau_i^+$$

Hermele, F

$$H_{\text{ring}} = - \sum_{\hexagon_p} \frac{K}{2} (\tau_1^+ \tau_2^- \tau_3^+ \tau_4^- \tau_5^+ \tau_6^- + h.c.),$$

$$H_{\text{LGT}} = \sum_{\langle \mathbf{r}\mathbf{r}' \rangle} \frac{U}{2} (E_{\mathbf{r}\mathbf{r}'} - \frac{\epsilon_{\mathbf{r}}}{2})^2 - \sum_{\hexagon_d} K \cos(\text{curl } A),$$

$H_{\text{LGT}}$  captures the **universal properties** of

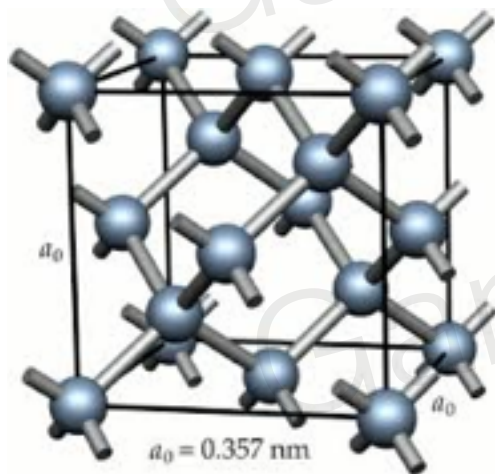
in the lattice gauge theory formulation, the  $S_z$  component is mapped to the lattice electric field, while  $S^+$ ,  $S^-$  is mapped to vector gauge potential.

the system is described by a lattice gauge theory on the diamond lattice formed by the center of the pyrochlore

in an ordered state,  $\langle S_z \rangle \neq 0$  from the Heisenberg relation,  $S^+$  is strongly fluctuating

In the gauge language, in an ordered state,  $E$  is static,  $B$  is strongly fluctuating, the magnetic monopole is condensed, which confines the electric charge.

the background monopole condensate disrupts the free motion of the spinons, the spinons are confined.



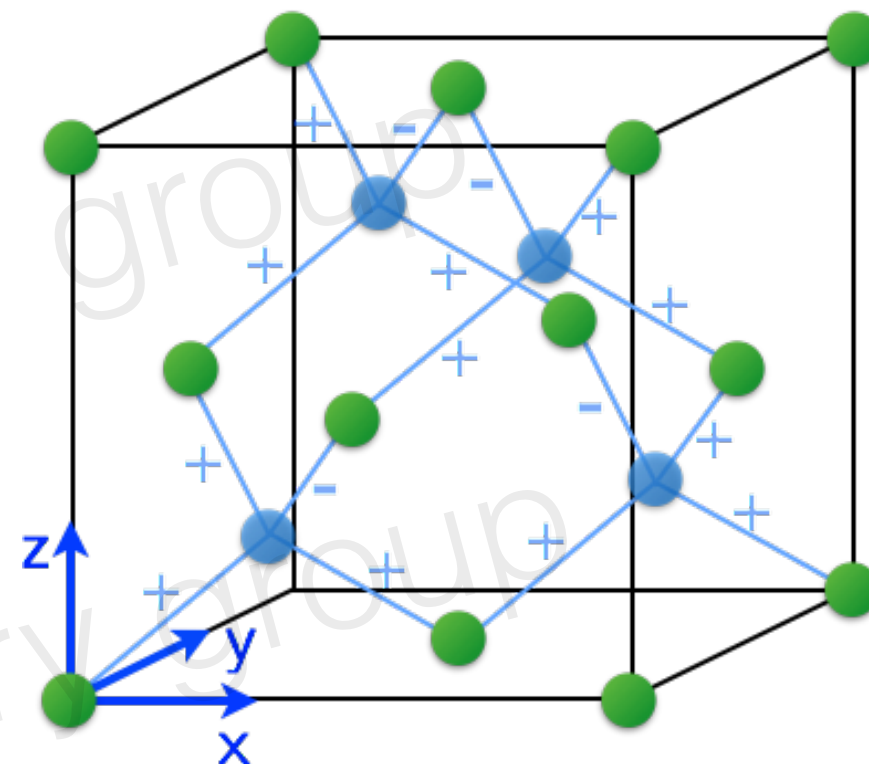
diamond lattice

- In an ordered state,  $\langle \tau_z \rangle \neq 0$ ,  $\langle \tau^+ \rangle$  is strongly fluctuating.
- In the gauge language, “**E field**” is static, “**B magnetic field**” is strongly fluctuating, the magnetic monopole (carrying magnetic charge) is condensed, which confines the electric charge carriers (spinons).

# Electromagnetic duality

Monopole lives on dual diamond lattice, carry magnetic charge or dual U(1) gauge charge.

$$H_{\text{dual}} = \sum_{\hexagon_d^*} \frac{U}{2} (\text{curl } a - \bar{E})^2 - \sum_{\mathbf{r}, \mathbf{r}'} K \cos B_{\mathbf{r}\mathbf{r}'} - \sum_{\mathbf{r}, \mathbf{r}'} t \cos(\theta_{\mathbf{r}} - \theta_{\mathbf{r}'} + 2\pi a_{\mathbf{r}\mathbf{r}'}).$$



monopole hopping on  
dual lattice

Motrunich, Senthil 2005,  
Bergman, Fiete, Balents 2006

Proximate magnetic order generically  
breaks translation symmetry

this is a bit technical.

I will explain by analogy

# Implication for experiments

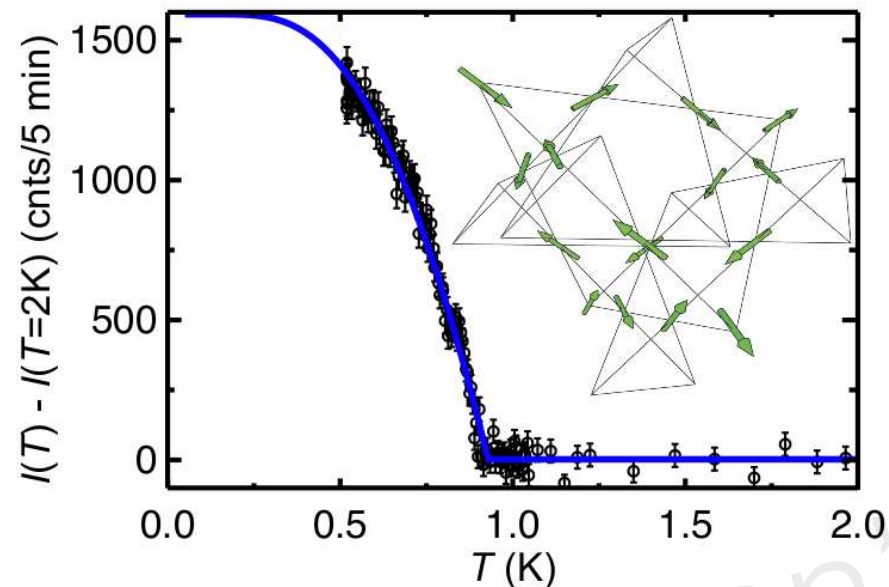


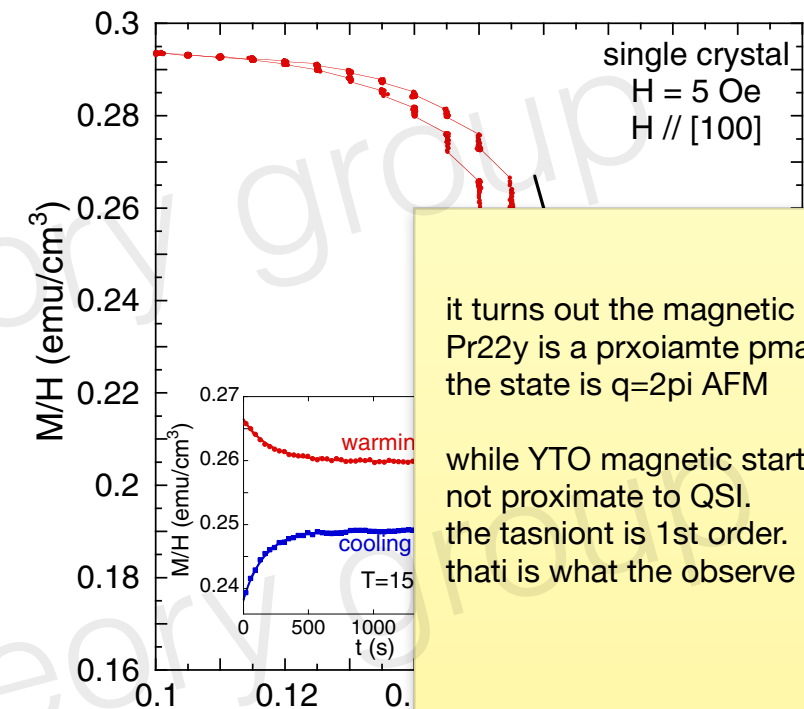
FIG. 2. (color online) Temperature dependence of elastic neutron scattering intensity of  $\text{Pr}_{2+x}\text{Ir}_{2-x}\text{O}_{7-\delta}$  at the position of the  $\mathbf{q}_m = (100)$  reflection. The intensity measured at  $T = 2$  K was subtracted as a background. Curve: Ising mean-field theory fit to the data, which yields a transition temperature of  $T_M = 0.93(1)$  K. Inset: sketch of the 2-in/2-out magnetic structure.

Magnetic order is discovered in some samples. (MacLaughlin, etc, 2015)

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it turns out the magnetic phase is Pr22y is a prxoiamte pmagnetic order, the state is  $q=2\pi$  AFM

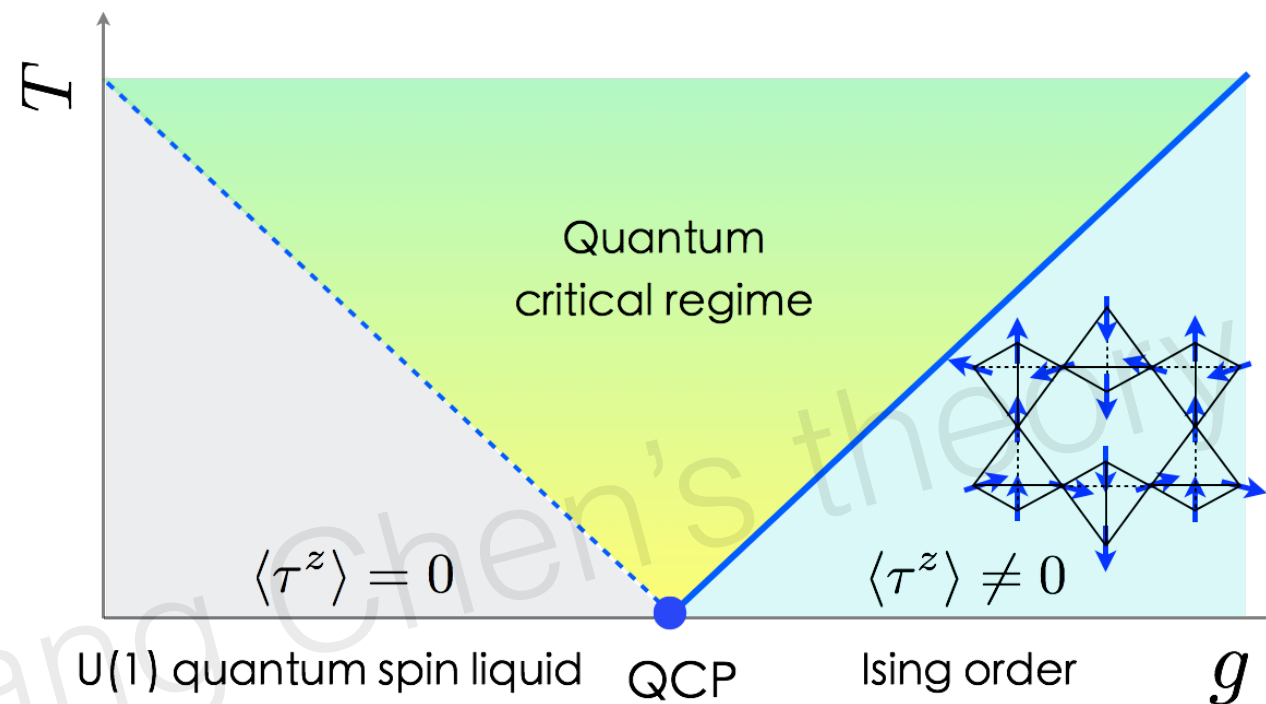
while YTO magnetic start is  $Q=FM$  is not proximate to QSI. the tasnion is 1st order. thati is what the observe in experiments

$\text{Yb}_2$

PIO: different samples have different Fermi energy  $\rightarrow$  RKKY- $\rightarrow$  **magnetic order,  $Q=2\pi(001)$**

YTO: First order transition to  **$Q=0$  FM state.**

# Subsidiary order and weak divergence



$g$  is the mass of the monopole

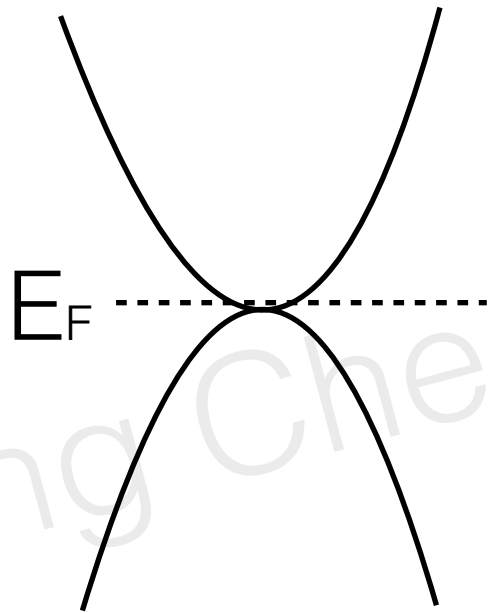
$$L = \sum_a \left[ |(\partial_\mu - i\tilde{a}_\mu)\phi_a|^2 + m^2|\phi_a|^2 \right] + \frac{F_{\mu\nu}^2}{2} + u_0 \left( \sum_a |\phi_a|^2 \right)^2 + u_1 \sum_{a \neq b} |\phi_a|^2 |\phi_b|^2 + \dots,$$

The critical theory is described by gapless monopoles coupled with a fluctuating U(1) gauge field in 3+1D.

a unusual weak divergence  $\chi(Q) \sim -\ln T$  “subsidiary order” (Kivelson) !



# More experimental prediction for $\text{Pr}_2\text{Ir}_2\text{O}_{7-\delta}$



Particle-hole excitations are centered at **Gamma** point

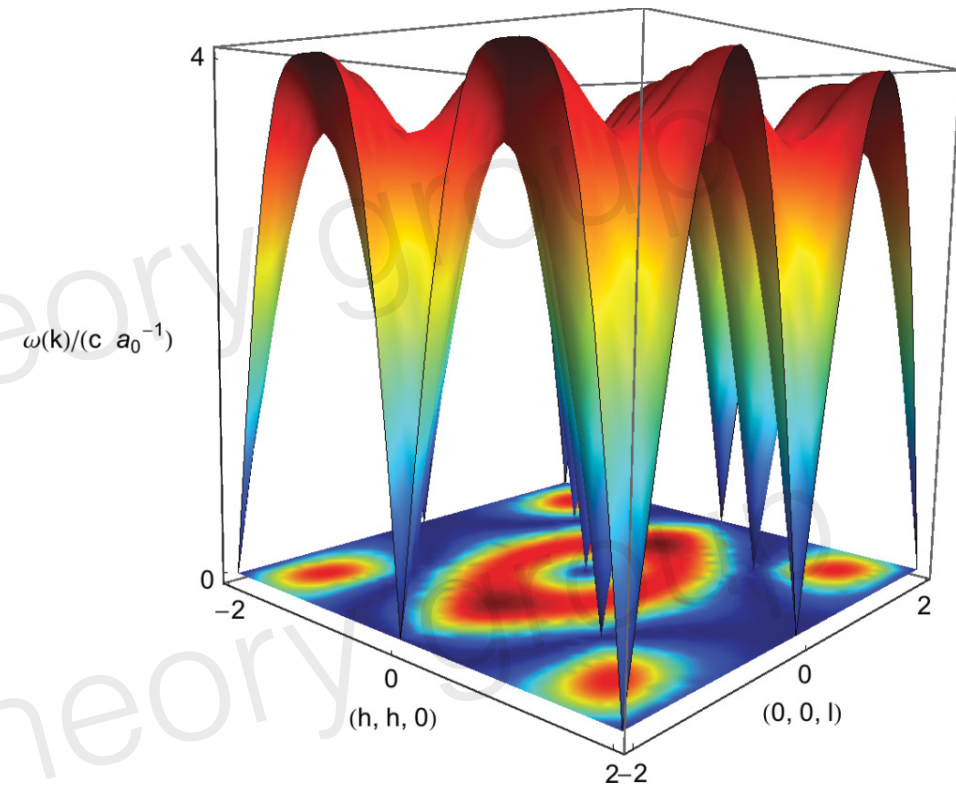


Fig from Benton, etc, PRB 2012

Emergent gauge photons are near the **suppressed pinch points**

The energy scales are different, maybe inelastic neutron scattering can work.

# Summary

- I have studied the phase diagram near quantum spin ice quantum spin liquid.
- Using field theoretic technique, I have obtained the structure of the magnetic states and the nature of the magnetic transition.
- I use the theoretical results to explain the puzzling experiments in  $\text{Pr}_2\text{Ir}_2\text{O}_7$  and  $\text{Yb}_2\text{Ti}_2\text{O}_7$ . It implies the disordered phase is a quantum spin ice U(1) quantum spin liquid.

**Ref:** Gang Chen, arXiv:1602.02230, longer talk can be found at KITP website last Sep.

**Work in progress:** sign problem free model that demonstrates both proximate and unproximate magnetic transition out of QSI QSL.

In preparation: Yao-Dong Li and Gang Chen, conduction electrons.